
Novel erbium-doped crystal saturable absorber Q-switches for the Er:glass laser

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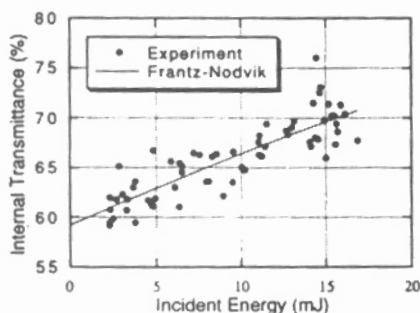
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Er:glass lasers are useful for a large number of applications because they emit in a narrow eye-safe region around $1.53 \mu\text{m}$. For many of these applications a Q-switched pulse is required. Active Q-switches increase the complexity and the cost of the system. Passive Q-switches such as saturable absorbers remain the simplest and cheapest method.¹⁻⁴

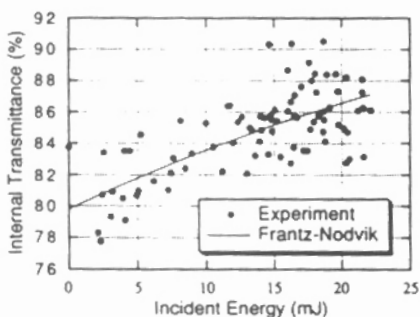
In this paper we report the first passive Q-switching of the Er^{3+} :glass laser at $1.53 \mu\text{m}$ using Er^{3+} : CaF_2 as an intracavity saturable absorber. Q-switched pulses of 69 ns (FWHM) and 11 mJ were obtained using an uncoated 1-mm-thick Q-switch sample in a 17.5-cm-long resonator cavity along with a $3 \times 50 \text{ mm}$ QE-7S (Kigre) Er:glass laser rod. Saturable absorber Q-switching of Er:glass has previously⁵ been obtained using Er: $\text{Ca}_5(\text{PO}_4)_3\text{F}$ (or Er:FAP); however, new results for this material, as well as a comparison to Er: CaF_2 are presented here. We achieved Q-switch pulses of 27 mJ, 47 ns using an AR-coated Er:FAP crystal, repre-

CTuK53 Table 1. Summary of Er:CaF₂ and Er:FAP Q-switch Results

Q-switch	Q-switch thickness (mm)	Er:glass laser rod	Resonator length/intracavity lens focal length	Output mirror reflectivity (%)	Output Energy (mJ)	Threshold (J)	Q-switch pulse width (nsec)
(3.5%)Er:CaF ₂	1.1	(QE-7) 3 x 76 mm	41 cm/+15 cm	88 (flat)	2.3	53	233
(3.5%)Er:CaF ₂	1.1	(QE-7S) 3 x 50 mm	21 cm/+7.6 cm	88 (flat)	3.3	205	129
(2.0%)Er:CaF ₂	1.0	(QE-7S) 3 x 50 mm	17.5 cm/+5 cm	94 (2.5 cm curv.)	11	51	69
(2.0%)Er:CaF ₂	1.0	(QE-7S) 3 x 50 mm	21 cm/+5 & -2.5 cm (intracav. telescope)	94 (flat)	1.6	39	400
(1%)Er:(1%)Yb:FAP(uncoated)	4.0	(QE-7S) 3 x 50 mm	17.5 cm/+5 cm	80 (2.5 cm curv.)	15	135	43
(1%)Er:(1%)Yb:FAP(coated)	4.5	(QE-7S) 3 x 50 mm	17.5 cm/+5 cm	94 (2.5 cm curv.)	27	44	47



CTuK53 Fig. 1. Bleaching data of (2%) Er:CaF₂ at 1.543 μ m. Solid curve is from modified Frantz-Nodvik equation for saturation fluence of 9.2 J/cm².



CTuK53 Fig. 2. Bleaching data of (1%)Er:(1%)Yb:FAP at 1.543 μ m. Solid curve is from modified Frantz-Nodvik equation for saturation fluence of 9.8 J/cm².

senting a significant improvement over the previously reported results (6 mJ, 83 ns).⁵

Our saturable absorber Q-switching results were obtained using several Er:glass resonator cavity configurations, summarized in Table 1. All the Q-switch crystals were cut and polished with flat, parallel surfaces. The Q-switches were aligned parallel to the cavity mirrors. The Er:FAP crystals (uniaxial) were cut with the c-axis parallel to the polished surfaces (E//c absorption band is higher and broader than E//a).

Q-switching with Er:FAP typically resulted in one very short pulse (Q-switched), followed by several free-running spikes of much lower peak power.⁵ The Er:CaF₂ Q-switches did not exhibit the free-running spikes but instead when the pumping energy was increased, we observed what seemed to be repetitive Q-switching.

The absorption cross-sections for both Er:CaF₂ and Er:FAP were determined by

bleaching the crystals using a Raman-shifted Nd:YAG laser at 1.543 μ m with a pulse-width of 15 ns. We measured 1.543 μ m saturation fluences (Figs. 1 and 2) of 9.2 J/cm² for Er:CaF₂ and 9.8 J/cm² for Er:FAP (E//c). These correspond to cross-sections of 1.41×10^{-20} cm² and 1.32×10^{-20} cm² (E//c) for Er:CaF₂ and Er:FAP, respectively. Using the absorption spectra for both materials, we corrected the cross-sections to 1.29×10^{-20} cm² and 1.43×10^{-20} cm² (E//c), respectively, for the Er:glass wavelength of 1.533 μ m. Theoretical modeling of the Q-switched pulses showed good agreement with the results summarized in Table 1.

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High-gain, diode-laser-pumped Nd:YVO₄ slab amplifier

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We have previously reported¹ a novel, diode-laser, side-pumped Nd:YVO₄ slab laser oscillator with high optical to optical conversion efficiency. This oscillator design uses a laser material such as Nd:YVO₄ which absorbs strongly at the pump wavelength so the laser gain is confined to a thin region next to the flat pump face. Good spatial overlap between the pumped volume and the laser mode is obtained by arranging the cavity optics so that the laser mode makes a single, grazing-incidence, total internal reflection from the pump face and therefore remains in the region of highest gain throughout its passage through the laser material. Because of the single reflection,