

Co<sup>2+</sup>:ZnSe Saturable Absorber Q-switch for the 1.54  $\mu\text{m}$  Er<sup>3+</sup>:Yb<sup>3+</sup>:Glass Laser

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### Abstract

A potentially superior eye-safe laser Q-switch material has been demonstrated. Pulses of 77 ns and 2.6 mJ were obtained in a non-focused Er:glass resonator.

### Key words

Transition metal doped materials, Infrared and far-infrared lasers, Optical properties of materials, Fluorescence.

### Introduction

We have demonstrated a new application for the important class of materials: transition-metal-ion-doped II-VI semiconductors as saturable absorber Q-switches.

Those semiconductors are very attractive because of their optical properties and their utility in laser devices. Co:ZnSe is a zinc-chalcogenide belonging to this class of materials. It is grown by modified vertical Bridgman process, crystallizing in a sphalerite or zinc-blend structure. Its lattice can be cubic or hexagonal with Co substituting Zn<sup>2+</sup>, in tetrahedral sites.<sup>1,3</sup> The tetrahedral sites are noncentrosymmetric and the even-parity components of the crystal field are weak, resulting in mid-infrared transitions.<sup>4</sup>

For Co<sup>2+</sup>:ZnSe, the d-d transitions are electric-dipole forbidden according to Laporte selection rules and as a consequence, cobalt radiative lifetime is quite long.<sup>1</sup>

Co<sup>2+</sup> has the electronic structure [Ar-core]3d<sup>7</sup>. Its outermost 3d-shell interacts strongly with lattice phonons, resulting in broad absorption bands. The first energy levels for a 3d<sup>7</sup>

configuration in tetrahedral symmetry are <sup>4</sup>A<sub>2</sub>(<sup>4</sup>F) or ground-state and the excited states <sup>4</sup>T<sub>1</sub>(<sup>4</sup>F), <sup>4</sup>T<sub>2</sub>(<sup>4</sup>F), absorbing from 1,200 to 2,000 nm and <sup>4</sup>T<sub>1</sub>(<sup>4</sup>P), absorbing from 680 to 800 nm. We are particularly interested in the broad absorption feature peaking at 1.6  $\mu\text{m}$ , covering the eye-safe range, from 1.5 to 1.8  $\mu\text{m}$ .<sup>3,5,6</sup>

Co<sup>2+</sup> in YSGG and in YAG were previously demonstrated by our group to function as saturable absorber Q-switches for the Er:glass laser.<sup>5</sup> However, the Co<sup>2+</sup>:garnet Q-switches required the use of an intracavity focusing lens.

It is advantageous to develop an efficient saturable absorber Q-switch with improved performance and reduced vulnerability to component damage, which does not require any intracavity focusing lens. The results reported in this paper indicate that Co<sup>2+</sup>:ZnSe may be the saturable absorber Q-switch of choice for the Er:glass laser.

In the Co:garnet crystals, the Co<sup>2+</sup> excited-state lifetime is very short (less than 1 ns).<sup>6</sup> For a fast relaxing Q-switch, one must consider the saturation intensity, which is proportional to  $(\sigma\tau)^{-1}$  where  $\sigma$  is the absorption cross-section, and  $\tau$  is the lifetime of the relaxation. Therefore, even materials which possess relatively high cross-section values can have high saturation intensities due to a short lifetime. The saturation intensities for Co in garnets is on the order of 100 MW/cm<sup>2</sup>.<sup>5</sup>

Co<sup>2+</sup>:ZnSe possesses a much longer excited-state lifetime ( $\sim 290 \mu\text{s}$ ), and a cross-section of  $\approx 7 \times 10^{-19} \text{ cm}^2$  at 1.53  $\mu\text{m}$ .<sup>1</sup> Thus, this material acts as slowly relaxing absorber, so that its high cross-section (about an order of magnitude higher<sup>7</sup> than U:CaF<sub>2</sub>) can be fully utilized. These considerations indicate that Co:ZnSe as a passive Q-switch material may perform even better than the U<sup>4+</sup>-doped difluoride crystals.<sup>7,8</sup>

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## Experimental Results

### Q-switch

A Q-switch sample was fabricated from a nominal 0.001%  $\text{Co}^{2+}:\text{ZnSe}$  boule fragment obtained from the Eagle-Picher Co.<sup>9</sup> We are interested in the  $\text{Co}^{2+}$  absorption band peaking at around 1.54  $\mu\text{m}$ , which is shown in Fig. 1. A sample thickness of 1.53 mm provided an internal loss of about 6%. Utilizing the cross-section data of reference 3, we estimate that the Co concentration in our sample was  $9 \times 10^{17} \text{ cm}^{-3}$ .

The experimental arrangement is shown in Fig. 2. The laser consisted of a flashlamp pumped Er:glass laser rod (3 x 50 mm), Kigre Er:Yb:phosphate glass, QE-7S, with an overall cavity length of 19 cm. The reflectivities of mirrors  $M_1$  and  $M_2$  were 100%, 50 cm radius of curvature, and 95%, flat respectively.

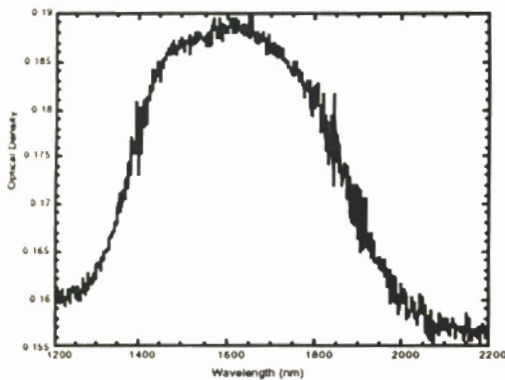


Figure 1. (0.001%) $\text{Co}^{2+}:\text{ZnSe}$ , 1.53 mm thick crystal, absorption spectrum at 1.54  $\mu\text{m}$ .

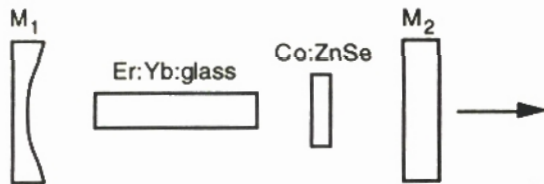


Figure 2. Experimental resonator cavity.

A typical output pulse is shown in Fig. 3. The full-width-half-maximum (FWHM) for the pulse is about 100 ns. The tail in the pulse is an artifact of the Judson InGaAs photodiode used in the experiment. In a shorter laser resonator (14.5

cm long), with a flat outcoupler of  $R = 85\%$ , we observed pulses of about 2.6 mJ and 77 ns.

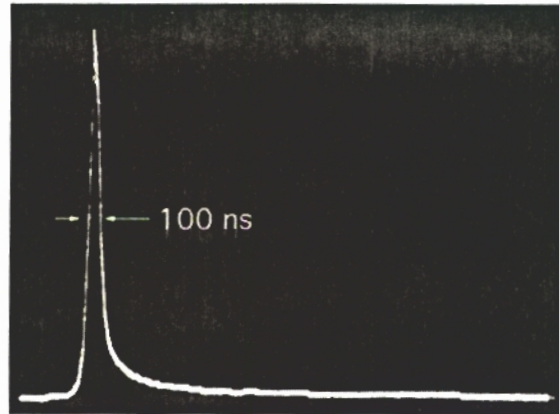


Figure 3. Q-switched pulse. Tail on pulse is an artifact of the detector used.

### Bleaching

The bleaching experiment was performed using a 1.4 mm thick, (0.1%) $\text{Co}^{2+}:\text{ZnSe}$  crystal and a Q-switched Er:glass laser (28 cm long laser resonator) at 1.53  $\mu\text{m}$ . The output energy was 10.1 mJ and the pulsewidth 60 ns. The laser beam was then focused over the sample with a +12.5 cm lens. The spot size at the sample was about 0.096 mm diameter.

We analyzed the bleaching results using a modified Avizonis-Grotbeck equation,<sup>10</sup> which includes excited state absorption (ESA):

$$\frac{dE}{dz} = -hvN_0 \left\{ \left( 1 - \frac{\sigma_2}{\sigma_1} \right) \times \left[ 1 - \exp\left( -\frac{\sigma_1 E}{hv} \right) \right] + \frac{\sigma_2 E}{hv} \right\} - \alpha E \quad (1)$$

where,  $E$  is the 1.53  $\mu\text{m}$  fluence and  $N_0$  is the  $\text{Co}^{2+}$  concentration.  $\sigma_1$  and  $\sigma_2$  are the ground-state and ESA cross sections, and  $\alpha$  corresponds to all the nonsaturable losses. Equation (1) was solved numerically, and the parameters  $\sigma_1$  and  $\sigma_2$  were adjusted to match the theory with the experimental data. We assumed  $\alpha = 0$ .

The plot in Fig. 4 represents the curve fitting with a contrast 40:1 (i.e.,  $\sigma_1/\sigma_2$ ). Hence, according to our preliminary measurements the effects of ESA are negligible. The value adjusted for  $\sigma_1$  was  $6.19 \times 10^{-19} \text{ cm}^2$  and the saturation fluence was 210  $\text{mJ}/\text{cm}^2$ .  $\sigma_2$  was on the order of  $1.55 \times 10^{-20} \text{ cm}^2$ .

Since the Q-switched pulse used to bleach the sample (60 ns) is much shorter than the

relaxation lifetime (290  $\mu$ s) of the absorber (i.e., slowly-relaxing regime), it is appropriate to discuss the saturation in terms of energy/cm<sup>2</sup>, or fluence. However, as a comparison to the 100 MW/cm<sup>2</sup> saturation intensity for the fast-relaxing Co: garnets, it is useful to note that the intensity corresponding to saturation fluence of Co:ZnSe is only 3.5 MW/cm<sup>2</sup>, in the bleaching experiment described above.

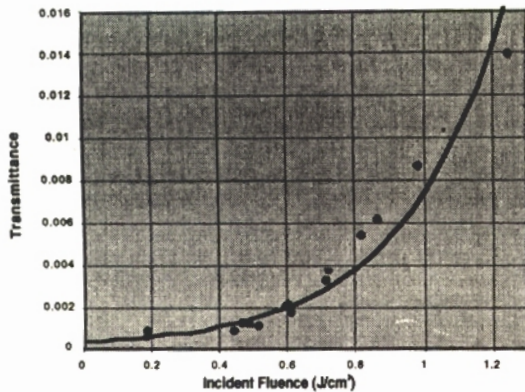


Figure 4: Transmittance as function of the incident fluence for a 1.4 mm thick (0.1%)Co<sup>2+</sup>:ZnSe crystal. The solid line represents the Avizonis-Grotbeck model.

### Theoretical Modeling

The rate equations for a slowly relaxing saturable absorber Q-switch with a three-level laser are given by:<sup>11</sup>

$$\frac{d\phi}{dt} = (K_g N_g - K_a N_a - \gamma_c) \phi \quad (2)$$

$$\frac{dN_g}{dt} = -\gamma K_g N_g \phi \quad (3)$$

$$\frac{dN_a}{dt} = -K_a N_a \phi \quad (4)$$

where  $\phi$  is the number of photons traveling in the cavity,  $\gamma_c$  is the photon-decay rate excluding the saturable loss, and  $N_g$  and  $N_a$  are the gain and the absorber media population differences.  $\gamma = 1 + g_2/g_1$ , where  $g_2$  and  $g_1$  are the upper and lower laser levels degeneracies for the laser media, respectively. We assumed  $g_1 = g_2$ , or  $\gamma = 2$ , for the Er:glass laser. The coupling coefficients  $K_i$  are given by  $K_i = \sigma_i/t_i A_i$ , where  $\sigma_{g,a}$  are the gain and

the absorber media cross sections and  $t_i$  is the one-way photon-cavity transit time.

Using the resonator and Q-switch parameters from the experiment (19 cm resonator,  $R = 95\%$ ) and solving numerically the saturable absorber Q-switch rate equations (2) through (4), pulses of 3.7 mJ and 58 ns are predicted. Losses due to the uncoated Co:ZnSe surfaces were included, and it was assumed that the entire 6% internal loss was saturable. The loss associated with the uncoated surfaces of the Q-switch can be substantial, and it is expected that AR-coatings, as well as an increase in the saturable losses (thicker Q-switch) will significantly shorten the pulsewidth and increase the output energy.

### Conclusion

We have demonstrated a new passive Q-switch material for eyesafe lasers, particularly the Er:glass laser, which preliminary results indicate to be superior to the present state-of-the-art (U:CaF<sub>2</sub>) saturable absorber Q-switches.

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