

**A SURVEY OF THORIUM UTILIZATION IN THERMAL POWER REACTORS**

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## I INTRODUCTION

Brazil has a natural interest in the thorium fuel cycle for nuclear reactors as it has large proven reserves of thorium and only small uranium deposits. In the thorium cycle, fissile U 233 is formed from Th 232. Thorium has been used in the cores of several power reactors in the mid sixties (Indian Point<sup>(1,2)</sup>, 250 MWe, Elk River 58 MWTh<sup>(3)</sup>, Borax IV). The experimental high temperature reactors Peach Bottom (40 MWe)<sup>(4,5)</sup>, AVR (15 MWe)<sup>(6)</sup> and Dragon (70 MWe)<sup>(7,8,9)</sup> are thorium reactors and also the 300 MWe prototype HTGR's of Fort St. Vrain and Schmehausen<sup>(10,11)</sup> will use the thorium cycle. A number of commercial size 1200 MWe thorium HTGR's have been ordered from General Atomic<sup>(12)</sup>.

Most work on the use of thorium in light and heavy water reactors has been discontinued because the low prices for uranium ore and separative work a number of years ago did not give any incentive to develop the thorium cycle further. Due to the ever increasing prices of uranium, the situation has changed now internationally and a renewed interest in the utilization of thorium can be observed.

Work on the utilization of thorium in high temperature reactors has been carried to the commercial stage. This is due to the fact that uranium is not competitive in homogeneous reactors like the HTGR. The uranium resonances in a homogeneous reactor are poorly shielded, leading to a low resonance escape probability. This means that a high enrichment is required to achieve criticality. The resonances of thorium are not so pronounced and the requirements of fissile material are thus less.

Water cooled reactors have much lower effective resonance integrals caused by the selfshielding of the heterogeneous fuel. The resonance integrals of thorium and uranium are for this reactor type about the same. Thorium has a disadvantage as its thermal capture cross section is about three times the value for uranium leading to a significant lower thermal utilization and thus requiring a higher fissile loading.

Thus far fuel cycles have been mainly compared by the fuel cycle costs. The cost of all steps of the fuel cycle mining, enrichment, fabrication and reprocessing were mainly those at the time the studies were made and the optimizations were accordingly<sup>(13)</sup>. Few consider the longterm trends in the costs and only the work on fast breeders has given attention to the availability of resources<sup>(14)</sup>.

The incentive to use thorium is its potential to conserve the resources. Uranium 233 that is bred from U Thorium is potentially a better fuel in thermal reactors than Pu-239 that is bred from U 238. U 233 does not occur naturally and has to be converted in thermal or fast reactors. Uranium has to be used to start up and maintain non breeding reactors as U 235 is the only

naturally occurring fissile material. It is thus not obvious if thorium utilization will indeed reduce the demand for uranium.

The requirements of uranium are mainly given by the conversion ratio of the reactors in the power system and depend on the layout of the reactor. In this paper we will show that with more emphasis on the conversion ratio, a substantial reduction of the ore requirements can be obtained and we will outline under what conditions thorium utilization will reduce the uranium consumption.

## 2 REACTOR TYPES

### 2.1 High Temperature Reactors

Thorium utilization in high temperature gas cooled reactors has reached a commercial stage and has been well documented. A comparison of the fuel cycle with uranium and thorium has been made by GA for prismatic fuel elements<sup>(15)</sup> and for pebbles (Shulter elements)<sup>(16)</sup> by Jülich. From the work of Teuchert we have taken Table I. This shows that the use of thorium gives only a limited reduction on the uranium consumption.

Table I

Characteristics of the fuel cycle

$1000\text{MW}_{\text{th}}$ ,  $8 \text{ MW/m}^3$ ,  $250-1050^\circ\text{C}$

Fuel cycle		Low. enr.	Thorium	Th. recycle	Hi. conv.
$N_C/N_{HM}$ in feed batches		355	259	259	110
Avg. fuel residence time	days	848	1094	1113	630
Average burn up	MWd/t	115 000	114 000	115 000	28 000
Conversion ratio		0.54	0.55	0.60	0.90
Power peaking		2.8	2.6	2.5	1.3
Control Rods Withdrawal	$\Delta k/k$	0.020	0.018	0.018	-0.006
Xe-Override 100 - 40%	$\Delta k/k$	0.016	0.019	0.018	-0.007
Fissile inventory	kg	345	361	360	620
supply	g/d	916	783	729	1082
U-nat. requirement	kg/d	217	190	135	-
Separative work	kg SWU/a	180	189	120	-
Fuel cycle costs	mills/KWh <sub>e</sub>	2.12	1.92	1.84	2.91
Fabrication only	mills/KWh <sub>e</sub>	0.49	0.41	0.41	1.20

## 2.2 - Light Water Reactors

There exist a relative large amount of information on the use of thorium in light water reactors from the prototypes in the early sixties. This information is however not too pertinent now as the cladding of the fuel elements used in these reactors consisted of stainless steel. Stainless steel is not now used as it has a high absorption cross section. The favorite material now is Zircaloy. A zirconium alloy with an extreme low capture cross section. The more recent work of Zorzi<sup>(17, 18, 19)</sup> focusses on the use of Th metals as a replacement for uranium oxide and stresses fuel cycle costs. A low fuel cycle cost is obtained and the uranium demand is significantly reduced. Lin and Zolotar<sup>(20)</sup> have done some zero dimensional burnup calculations comparing  $UO_2$ ,  $ThO_2$  and Th metal using the code LEOPARD<sup>(21)</sup>. Their results are of a preliminary nature and are given in Table II.

Table II

PWR Fuel Cycle Characteristics

	$UO_2$ - $PuO_2$	$ThO_2$ - $UO_2$	Th-U
Initial uniform loading - MT $U^{235}$	2.740	3.681	4.583
Net U (MT) <sup>*</sup>	$6.028 \times 10^2$	$7.189 \times 10^2$	$8.951 \times 10^2$
Separative work (Kg) <sup>*</sup>	$4.063 \times 10^4$	$9.324 \times 10^5$	$11.161 \times 10^5$
Conversion ratio of first cycle	0.81	0.76	0.81
Make up requirement per year - Mt $U^{235}$	0.449	0.315	0.176
Net U (MT) <sup>*</sup>	$0.824 \times 10^2$	$0.615 \times 10^2$	$0.344 \times 10^2$
Separative work (Kg) <sup>*</sup>	$1.018 \times 10^{3**}$	$0.798 \times 10^2$	$0.446 \times 10^5$
Conversion ratio of second cycle	0.87	0.79	0.84

\* 0.2 w/o  $U-238$  in diffusion plant tails.

\*\* Includes enrichment of recycled uranium.

The light water breeder reactor program is intended to develop a thermal breeder reactor<sup>(22,23)</sup> on the basis of the existing light water reactor technology. A conversion ratio of nearly one is obtained which implies that except for the initial loading, practically no new fuel is needed. But the complexity of the seed and blanket concept and specifically the movable seed reactivity control make it very doubtful if a high availability can be obtained.

## 2.3 - Heavy Water Reactors

Some work has been performed on heavy water reactors to obtain breeding using the thorium cycle. For the suspension power reactor<sup>(24)</sup> a homogeneous fueled heavy water reactor a conversion ratio of 80 is pressurized heavy water reactor - a concept similar to the standard PWR but moderated and cooled with heavy water and larger distances between fuel pins. Conversion ratios close to 1.05 for thorium loadings have been obtained<sup>(25)</sup>. Work on a similar reactor type has been performed by Siemens in Germany<sup>(26)</sup>. The large pressure vessels for these reactors can possibly be made from prestressed concrete<sup>(27)</sup>. The natural uranium cycle for these reactors is not attractive as the burnup reach only 4 MWd/kg.

Thorium utilization in heavy water moderated organic cooled reactors of the pressure tube type CANDU ORGEL have been extensively studied<sup>(29 30 31 32 33)</sup>, and offers a high conversion ratio with low fuel cycle costs. One can start CANDU type reactors with natural uranium, reprocess the fuel and use the plutonium in thorium elements so as to start a U 233 Th cycle. But this way though promising, takes about 30 years to build the necessary inventories<sup>(28)</sup>.

#### 2.4 - Molten Salt Breeder Reactor

An advanced and interesting thermal converter is the molten salt breeder reactor<sup>(35 36)</sup>. It combines a low inventory with an high conversion ratio. It is likely that with this system thermal breeding will be possible. The MSBR consists of a graphite matrix in which a mixture of lithium, uranium and thorium fluorides are passed through. The high conversion ratio is obtained by online reprocessing of the fuel mixture, thus removing virtually all fission products. The reactor has no structural material and the moderator does not absorb neutrons. These factors explain the high conversion ratio. Some basic problems, it seems, require a lot of development work.

One is a corrosion resistant alloy for the reactor vessel and piping as the fission products in the fluoride solution can be aggressive.

A large amount of tritium is formed by the reaction  $\text{Li}^6(n,\alpha)\text{T}^3$  about 2.4 KCi/d, and unless suitable measures are taken, an excessive amount of the tritium could reach the atmosphere by diffusion through the heat exchangers into the steam system.

Solutions to these problems do exist, but they require extensive testing on an engineering scale. In the opinion of this author the development of the MSBR is lagging behind the liquid metal fast breeder reactor by about a decade and will probably not be in commercial operation before the next century.

### 3 PHYSICS CONSIDERATIONS

As an example of the effects that are involved in the transition from the uranium cycle to the thorium cycle we will have a more detailed look at the light water reactors. First approximations are used to demonstrate the effects that are involved. The results, though basically correct, should be treated with some care as second order effects can significantly modify them. This is particularly true for conversion ratios close to unity. Small changes in this number lead to large changes in the fuel consumption.

#### 3.1 - Conversion Ratio

The basic quantity governing the fuel consumption is the conversion ratio. In comparing the conversion ratio of thorium and uranium reactors one should be careful to separate the effects of U 233 and Pu 239 from other effects like change in enrichment, burnup, cladding absorption, absorption in moderator and control requirements.

The number of neutrons that are available for conversion per absorption in a fissionable atom are given by  $\eta - 1$  of that isotope. We have listed the  $\eta$  together with the fission cross

section and the ratio of absorption to fission for the most important isotopes in table III

Table III

Thermal Properties of Some Fissile Isotopes

	$\sigma_f$	$\eta$	$\alpha$
U 233	276	2.3	1
U 235	277	2.0	2
Pu 239	786	1.8	4
Pu 241	766	2.2	2

Figures 1 to 4 give graphs of some of these quantities

In the standard light water reactor, plutonium is built up during the irradiation and contributes about 30% to the fissions. The average  $\eta$  is thus  $\eta_{LWR} = 2$ . The fast fissions in uranium increase the number of neutrons available to breeding to  $\epsilon\eta - 1 = 1.1$ . This number is reduced to .6 by capture in cladding, coolant, fission products and control absorber. In replacing uranium 238 with Th 232, we will have several compensating effects. First of all the fast fission bonus is practically lost because Th 232 has a lower fast fission cross section, secondly because the  $\eta$  of U 233 is so much higher than that of Pu 239 we obtain for  $\epsilon\eta - 1 = 1.1$ . The reduction due to capture in the intermediate isotope Pa 233 is not very significant. We can thus conclude that the conversion ratio of a light water reactor without recycling is not affected by the substitution of U 238 by Th 232 for the same enrichment, burnup etc

### 3.2 Critical Mass

What one would do in an existing light water reactor is to substitute the  $UO_2$  in the fuel pins with a mixture of  $ThO_2$  and highly enriched  $UO_2$  as was done in the early light water reactors.  $ThO_2$  behaves very similar to  $UO_2$  so that this would not present significant problems

The absorption of thorium in standard LWR's is about 20% higher than that of U-238 which is caused by thermal absorption cross section that is three times higher. The effect of the thermal absorption is somewhat compensated by an effective resonance integral that is only 75% of that of U-238

The increase in absorption together with the loss of fast fissions means that the fissile mass has to be increased by 20% in order for the system to remain critical as the density of  $ThO_2$  is only 90% of  $UO_2$ . Th metal has a 20% higher density and, if that is used, the enrichment should be increased by 60%



This increase in enrichment necessary for Th metal results in a conversion ratio of 75 a value close to that obtained by Zorzoli<sup>(28)</sup> when his values are extrapolated to burnups of 30 MWd/kg. The same conversion ratio can be obtained for a uranium system by reducing the fuel to moderator ratio because the increase in conversion ratio is a result of a change in enrichment.

### 3.3 Recycling

We will now consider the effect of recycling the fuel. We will consider two types of reactors, one type we call a converter that not only produces electricity but converts fertile isotopes to fissile e.g. U 238 to Pu 239. The other type we call a burner that burns converted fissile material both of itself and of true converters.

The net production of bred fuel in a converter is

$$P_{\text{net}} = \left\{ (1 + \alpha_5) (1 - f) + (1 + \alpha_b) f \right\} CR_c - (1 + \alpha_b) f \Big/ X \quad (1)$$

where

- $\alpha_5$  indicates U 235
- $\alpha_b$  indicates bred material, e.g., Pu 239 or U 233
- $f$  is the fraction of the power generated in the bred fuel
- $X$  is the power per kg of fissioned material
- $CR_c$  is the conversion ratio

The first term is the total production and the second term is the loss of bred material by fission and absorption in the converter itself.

The net consumption of a burner, assuming that the bred material is recycled, is given by

$$C_{\text{net}} = (CR_b - 1) (1 + \alpha_b) / X \quad (2)$$

where

$CR_b$  is the conversion ratio of the burner

From these equations we can now calculate the relation between the number of converters and the number of burners.

In Table IV we have given the values for the thorium and uranium cycle. The standard case is a thorium converter with the same enrichment as the standard PWR. The high conversion case is a thorium metal fed standard PWR with an adjusted enrichment and the equivalent uranium reactor. As a variant on the U-Th metal reactor we have included a Pu-Th metal reactor.

The different conversion ratios of the burner reflect the distinct properties of U-235, U-233 and Pu-239. The large consumption of plutonium in the burner in the uranium cycle, is a striking fact. This is due to the low conversion ratio and to the high  $\alpha$  of plutonium. If we have

a power system with converters and burners in the relation as given in Table IV we can calculate the average fuel consumption per reactor as given in Table V

Table IV

## Production and Consumption of Reactor Fuel

	Standard	Standard	High Conversion		
	U cycle	Th cycle	U cycle	Th cycle	Pu Th cycle
$CR_c$	6	6	75	75	65
$CR_b$	.45	75	60	90	90
$P_{net}$ (kg/y)	160	240	300	360	480
$C_{net}$ (kg/y)	800	220	520	90	90
Ratio of converter:					
Burner	5.1	1.1	3.2	1.4	1.5

Table V

Average Fuel Consumption per Reactor. \*Pu consumption \*\*\$30/kg  $U_3O_8$  and \$47/SWU

	Standard	Standard	High Conversion		
	U cycle	U cycle	U cycle	Th cycle	Pu Th cycle
U 236	500	320	400	130	670*
$U_3O_8$ ( $10^3$ kg/y)	100	64	80	25	
Enrichment (SWU)	66	100	50	40	
Costs** ( $10^6$ \$/y)	6.	5.8	4.8	2.4	

It can be observed that the reduction in uranium consumption is substantial, for the thorium cycle. This is offset by the increased enrichment costs as it is twice expensive to obtain 93% enriched material than 3% enriched.

For the high conversion system a factor of four reduction is achieved in ore consumption and nearly a factor of two in separative work requirements.

It is possible to start the thorium cycle with plutonium as a fuel and a nice relationship between Pu thorium and U 233 thorium reactors is obtained. It takes however, 1500 kg or ten/discharges of present day light water reactors for the loading of a Pu-thorium converter, and eight discharges of this converter are needed to load a U-233-thorium burner.

It is thus evident that long time periods are involved in building up U-233-thorium burners using plutonium as an intermediate fuel.

### 3.4 The use of 20% enriched uranium

Thus far we have assumed that 93% enriched uranium would be available. This might not be the case and the enrichment could be limited to 20%. There are two possibilities to build up the U-233 thorium cycle under these conditions.

One is to use plutonium as an intermediate step. To produce the 870 kg/y fissile plutonium that are needed for a Pu-Th metal reactor the output of five conventional light water reactors is needed. In equilibrium conditions one standard light water reactor produces enough fissile material to sustain one U-233 Th metal burner intermediate Pu converter. If highly enriched material is available, one LWR uranium converter could sustain four burners. This is a loss of a factor of four in ore consumption and a factor of two in separative work.

The other possibility is to use the uranium at the maximum permissible enrichment. Then, plutonium is produced together with U-233. The overall enrichment for a Th-metal reactor has to be of the order of 6%. Thus 30% of the heavy metal is U-238. The capture cross section of U-238 increases by a factor of two due to the dilution and consequent reduction in self-shielding of the resonances. Thus about 50% of the production of fissile material will be in the form of plutonium.

To obtain an idea of the mass balances involved we assume that the bred U-233 can be kept separate from the U-238 by a separation of fissile material. Then, we assume that we construct a system of ten light water reactors as converters, twenty thorium burners supplied with the U-233 fuel from the converters, one plutonium converter supplied with the plutonium of the ten uranium converter and five thorium burners fueled from the plutonium converter. Thus one light water reactor can sustain 3.5 thorium burners.

It is obvious that heavy penalties are involved should the uranium enrichment be limited to 20%. It probably does mean that the thorium cycle can not economically be developed.

## 4 REPROCESSING AND REFABRICATION

It has been observed<sup>(37)</sup> that the lack of a large scale facility for thorium reprocessing has been a principal barrier to the development of the thorium cycle.

Reprocessing of thorium based fuels can be done at existing facilities optimized for uranium and plutonium recovery but the cost is much higher. A plant optimized for thorium recovery can be made competitive to those of uranium plutonium<sup>(38)</sup>. Work is in progress in the United States and Germany to come to large scale reprocessing of HTGR fuel. The German Jupiter plant will use, after a head end that removes the bulk of the graphite, a Thorex solvent (13 M HNO<sub>3</sub>, 1M Al(NO<sub>3</sub>)<sub>3</sub>, 0.5 M HF) and extracts uranium, thorium and plutonium isotopes from this solvent with TBP<sup>(39)</sup> (Fig 5). This is a modified form of the well known pyrex process. In the U.S. a similar scheme will be used<sup>(40)</sup>. These processes can be adapted for water reactor fuel by a different head end to remove the bulk of the cladding.

A problem for the fabrication of fuel elements containing recycle fuel is the radioactivity. In the uranium cycle this is due to the  $\alpha$  activity of the plutonium isotopes and glove box handling is considered to be sufficient.

For the thorium cycle there are some added problems due to the buildup of U 232 (Fig. 6) The decay chain (Fig. 7) involves the production of  $\gamma$  rays with energies over one MeV that require heavy shielding (Table VI). Feed U 232 decays to Th 228 with a half-life of 2 y. Its daughters have short half-lives ( $< 1$  d). Thus, low radiation levels can be chemically obtained by removing Th 228 from the uranium feed before reprocessing is done. Th 228 cannot be separated from the thorium itself so that one has to wait a decade till the Th 228 has decayed before the thorium can be reused. Under these conditions glove box handling is sufficient and no refabrication penalty is incurred with respect to plutonium.

Experience has been obtained on a pilot plant scale by the Babcock and Wilcox Company on the fabrication of U 233 for water cooled reactors<sup>(41)</sup> as seen in Fig. 8.

The radioactive nature of the uranium 233 and the recycle thorium made it seem worthwhile to consider a close coupling between reprocessing and refabrication in which the used fuel is only partially decontaminated and the fabrication is done remotely.

Two such facilities have been constructed, the thorium uranium recycle facility (TURF) at Oak Ridge<sup>(3)</sup> and "programma ciclo uranium torio (PCUT)" in Italy. TURF has been designed mainly for HTGR elements whereas PCUT<sup>(42)</sup> was designed for light and heavy water fuel elements.

Essentially, in these facilities it is used the Sol-Gel process, is used for the oxide production which is then vibratory compacted in the fuel elements for water reactors. These steps are very much suited to remote handling (Fig. 9).

## 5 PROPERTIES OF THORIUM METAL AND THORIA

The properties of thorium, its alloys and compounds have been reported by Peterson et al.<sup>(43)</sup> Thoria behaves very much the same as uranium dioxide and does not pose particular problems (Table VII). Metallic thorium (Table VIII) is very attractive because it has a higher density than  $UO_2$  and also has a high thermal conductivity. Its irradiation behaviour is considered to be excellent and burnups of 30 MWd/kg should be obtained provided the maximum temperature is not much greater than 650°C<sup>(44)</sup>. The large swelling above 650°C (Figs. 10 & 11) is probably caused by the agglomeration of fission products in pores and it is not clear whether the swelling levels of 8%  $\Delta V/V$  per % burnup will be exceeded at temperature above 1000°C.

## 6 CONCLUSIONS

The thorium cycle offers substantial savings in uranium consumption and separative work for reactor with a high conversion ratio. The reactor that can be used are of a standard type and will need only slight modifications. A very interesting type is the Th metal pressurized water reactor into which existing pressurized water reactors can be converted.

Fuel fabrication facilities to be set up in Brasil should be able to handle slightly radioactive fuel and provide for removal of Th 228 from uranium.

Reprocessing plants should be optimized for a large through put of thorium even if initially the uranium cycle will be used.

A limit on the enrichment of feed uranium will at least double the uranium consumption.

Table VI

## Important Properties of Thorium and Uranium Dioxides

Property	Value for ThO <sub>2</sub>	reference	Value for UO <sub>2</sub> (24)
Crystal structure	Face-centred cubic (CaF <sub>2</sub> type)		Face-centred cubic (CaF <sub>2</sub> type)
Space group	O <sub>h</sub> <sup>5</sup> Fm3m		O <sub>h</sub> <sup>5</sup> Fm3m
Lattice parameter (Å)	5.5974 at 26°C 5.6448 at 942°C	(32)	5.4704 at 20°C 5.5246 at 946°C
Theoretical density (g/cm <sup>3</sup> )	10.00		10.96
Interatomic distances (Å)			
M-M	3.958		3.868
M-O	2.799		2.735
M-O	2.424		2.368
Thermal properties			
Melting point (°C)	3300 ± 100	(30)	2760 ± 30
Spectral emissivity (λ = 0.65 μm)	0.53 at 300°C to 0.21 at 800°C 0.2 to 0.65 at 1300°C depending on sample history	(33) (34)	0.416 ± 0.026 (near m.p.) 0.850 at 727°C 0.370 at 1847°C
Thermal conductivity (W/cm deg C)	0.103 at 100°C 0.034 at 800°C 0.086 at 200°C 0.031 at 1000°C 0.080 at 400°C 0.025 at 1200°C 0.044 at 600°C	(35)	0.106 at 100°C 0.0915 at 200°C 0.0560 at 400°C 0.0462 at 600°C 0.0376 at 800°C 0.0351 at 1000°C
Heat capacity (cal/mole deg C) (298 to 1200°C)	17.060 + 18.06 <sup>-4</sup> T - 2.5166(10 <sup>-5</sup> )/T <sup>2</sup>	(26)	18.45 + 2.431(10 <sup>-3</sup> )T - 2.272(10 <sup>-5</sup> )/T <sup>2</sup>
Debye temperature (°K)	200		870°K < 600°K (300-600°K)

Table VI (cont.)

Property	Value for ThO <sub>2</sub>	Ref.	Value for UO <sub>2</sub>
Coefficient of linear expansion ( $^{\circ}\text{C}^{-1}$ )	$62.16 (10^{-6} + 3.241 \cdot 10^{-9}) T - 0.1125/T^2$	36	$10.8 \times 10^{-6}$ (20-926 $^{\circ}\text{C}$ ) $9.8 \times 10^{-6}$ (75-800 $^{\circ}\text{C}$ ) $10.0 \times 10^{-6}$ (400-900 $^{\circ}\text{C}$ ) $10.52 \times 10^{-6}$ (26-1000 $^{\circ}\text{C}$ ) $\log P = 33.15/T - 4.026 \log T + 23.111$
Vapour pressure (atm)	$7.64 \cdot 3.3440(10^4)/T(2200-2900^{\circ}\text{K})$ $9.02 \cdot 3.78 (10^9)/T(2170-2400^{\circ}\text{K})$	37 38	
Thermodynamic properties			
Heat of formation, $\Delta_f H^{\circ}$ (kcal/mole)	$29.2 \pm 0.4$	29	$259.2 \pm 0.6$
Free energy of formation, $-\Delta_f G^{\circ}$ (kcal/mole)	$279.2 \pm 0.65$		$246.6 \pm 0.6$
Entropy, 298 $^{\circ}\text{K}$ , (cal/mole deg C)	$15.593 \pm 0.02$	26	$18.6 \pm 0.1$
Heat of sublimation (kcal/mole)	$158.7 \pm 2.5$ in range 2000-3000 $^{\circ}\text{K}$	37	$137.1 \pm 1.7$ at 1800 $^{\circ}\text{K}$
Entropy of sublimation (cal/mole deg C)	$35.3 \pm 1$ in range 2000-3000 $^{\circ}\text{K}$	37	$38.4$ at 1800 $^{\circ}\text{K}$
Mechanical properties			
Elastic properties			
Young's modulus (kilobars, lb/in <sup>2</sup> )	$1370, 19.8 (10^5)$	23	$1930, 26.0(10^5)$ at room temperature $1827, 26.5(10^5)$ at room temperature $1658, 24.0(10^5)$ at 800 $^{\circ}\text{C}$ $745, 10.8(10^5)$ at room temperature
Shear modulus (kilobars, lb/in <sup>2</sup> )	$990, 14, 3110^6$ at 30 $^{\circ}\text{C}$ $390.5, 8(10^6)$ at 1300 $^{\circ}\text{C}$	39	$0.302$ at room temperature $0.965, 1, 10, 14, 000 - 15, 000$ $4, 14-66, 60, 000 - 140, 000$ $0.16 - 0.37, 2, 300-5400$ $906 \pm 14$ 626
Poisson's ratio	0.17		
Modulus of rupture (kilobars, lb/in <sup>2</sup> )	0.83, 12, 000	40	
Compressive strength (kilobars, lb/in <sup>2</sup> )	15, 214, 000	23	
Fracture strength (kilobars, lb/in <sup>2</sup> )	1, 14, 000	23	
Hardness, Knoop	640 (500-g load)	41	
Electrical, magnetic, and optical properties			
Electrical resistivity (ohm cm)	$1$ to $> 10^4$ at 1600 $^{\circ}\text{C}$ , depending on sample and treatment	23	$3 \times 10^{-4} - 10^{-6}$ at room temperature about $10^{-1}$ at 500 $^{\circ}\text{C}$ about $10^1$ at 1000 $^{\circ}\text{C}$ 2.36 1.2 - 2.7
Index of refraction	2.09		

Table VII

## Physical and Mechanical Properties of Thorium Metal

Lattice parameter, A	
Face-centered cubic structure up to 1400°C (2550°F)	$\alpha_0 (25^\circ\text{C}, 77^\circ\text{F}) = 5.086 \pm 0.001$
Body-centered cubic structure 1400°C (2550°F) to melting point	$\alpha_0 (1450^\circ\text{C}, 2640^\circ\text{F}) = 4.11 \pm 0.01$
Density, g/cm <sup>3</sup>	
Theoretical	11.72
Calcium reduced (as cast)	11.5 - 11.6
Arc melted iodide	11.66
Melting point, °C (°F)	1690-1750 (3075-3180)
Boiling point, °C (°F)	3300-4500 (5975-8130)
Heat of vaporization, Kcal/mole	130-177
Heat of fusion, Kcal/mole	< 4.6
Elastic constants (25°C, 77°F)	
Modulus of elasticity, psi	$10.3 \times 10^6$
Shear modulus, psi	$4.1 \times 10^6$
Poisson's ratio	0.27
Compressibility, cm <sup>2</sup> /dyne	$16.4 \times 10^{-11}$
Thermal conductivity, 100-650°C (212-1200°F) cal/sec/cm <sup>2</sup> /°C/cm	0.090-0.108
Coefficient of linear expansion ( $\alpha_v$ ), °C <sup>-1</sup>	
25-200°C (77-390°F)	$11.0 \times 10^{-6}$
25-1000°C (77-1830°F)	$12.5 \times 10^{-6}$
Electrical resistivity, $\mu\text{ohm}\cdot\text{cm}$	
Calcium-reduced metal	18
Estimated for pure thorium	13-14
Self-diffusion coefficients, cm <sup>2</sup> /sec	
1100-1400°C (2010-2550°F)	$2 \times 10^{-10}$ - $2 \times 10^{-9}$
1450-1550°C (2640-2825°F)	$2 \times 10^{-8}$ - $1 \times 10^{-7}$

RADIATION ENERGIES OF THORIUM DECAY PRODUCTS<sup>1</sup>

Alpha decay			Beta decay		
Isotope	Mev	Mev	Isotope	Mev	Mev
U - 232	5.3	0.06, 0.29			
Th-232	4.0, 4.2		Ra-228(Ms Th <sub>1</sub> )	0.002	
Th-228(RdTh)	5.4	0.09	Ac-228(Ms Th <sub>2</sub> )	1.55	0.97
					0.91
Ra-224(ThX)	5.7, 5.4	0.25			0.46
					0.33
Rn-220(Th)	5.3				0.06
Po-216(Th A)	6.8				
Bi-212(Th C)	6.1	0.29, 0.14	Pb-212(Th B)	0.33(88%)	0.24(90%)
				0.57(12%)	0.30(9%)
					0.12(2%)
					0.18(4%)
					0.72(18%)
					0.83(16%)
					1.03(6%)
					1.35(6%)
					1.61(12%)
					1.80(10%)
			Tl-208(Th C <sup>''</sup> )	1.79	2.62(40%)
					0.58(40%)
					0.51(20%)

TABLE VIII



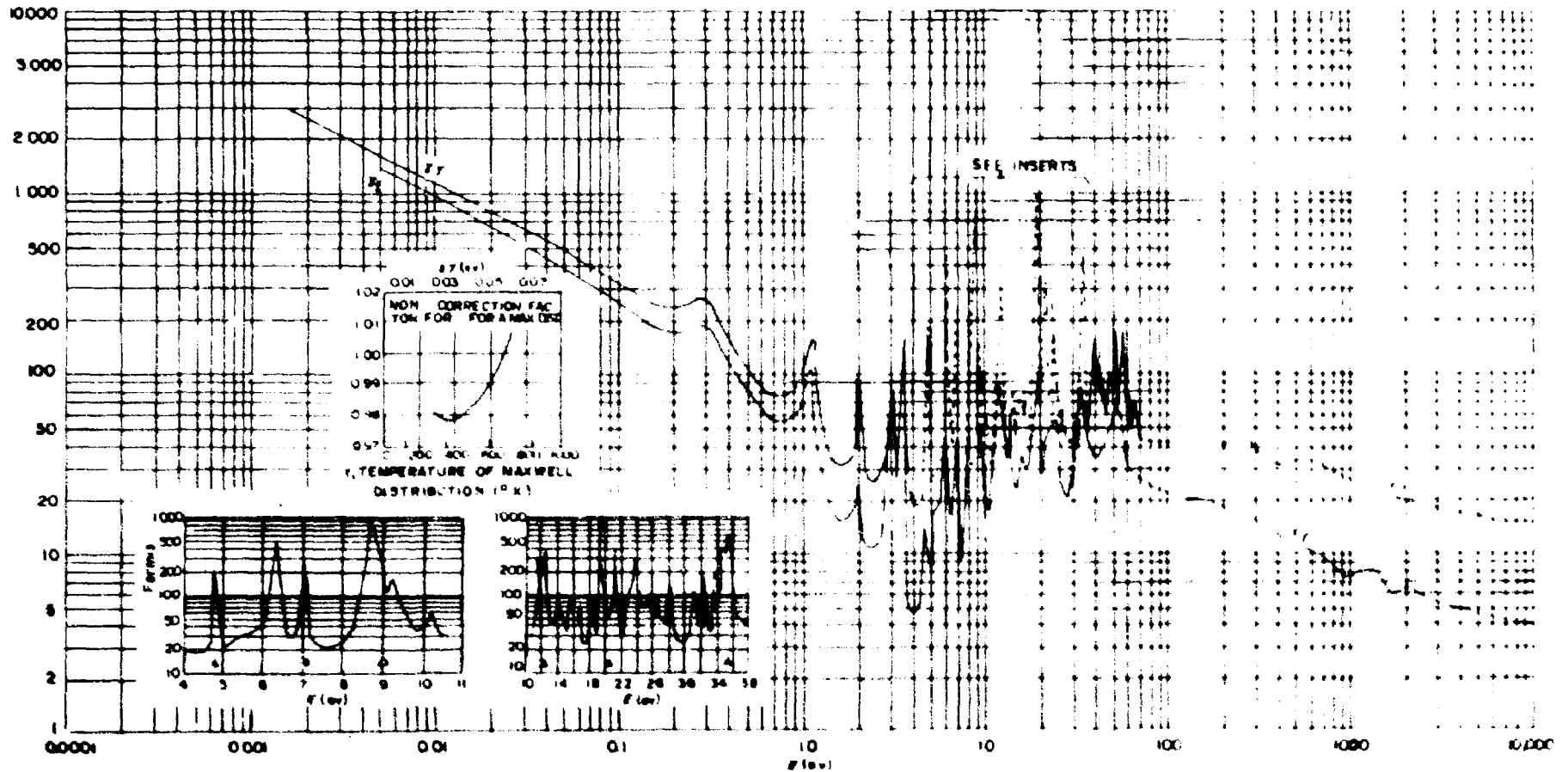


Fig. 1 - Total ( $\sigma_T$ ) and fission ( $\sigma_f$ ) cross-section of  $U^{235}$  in the low resonance and thermal regions. The triangles in  $\Delta E$  and the next three figures give the resolution width

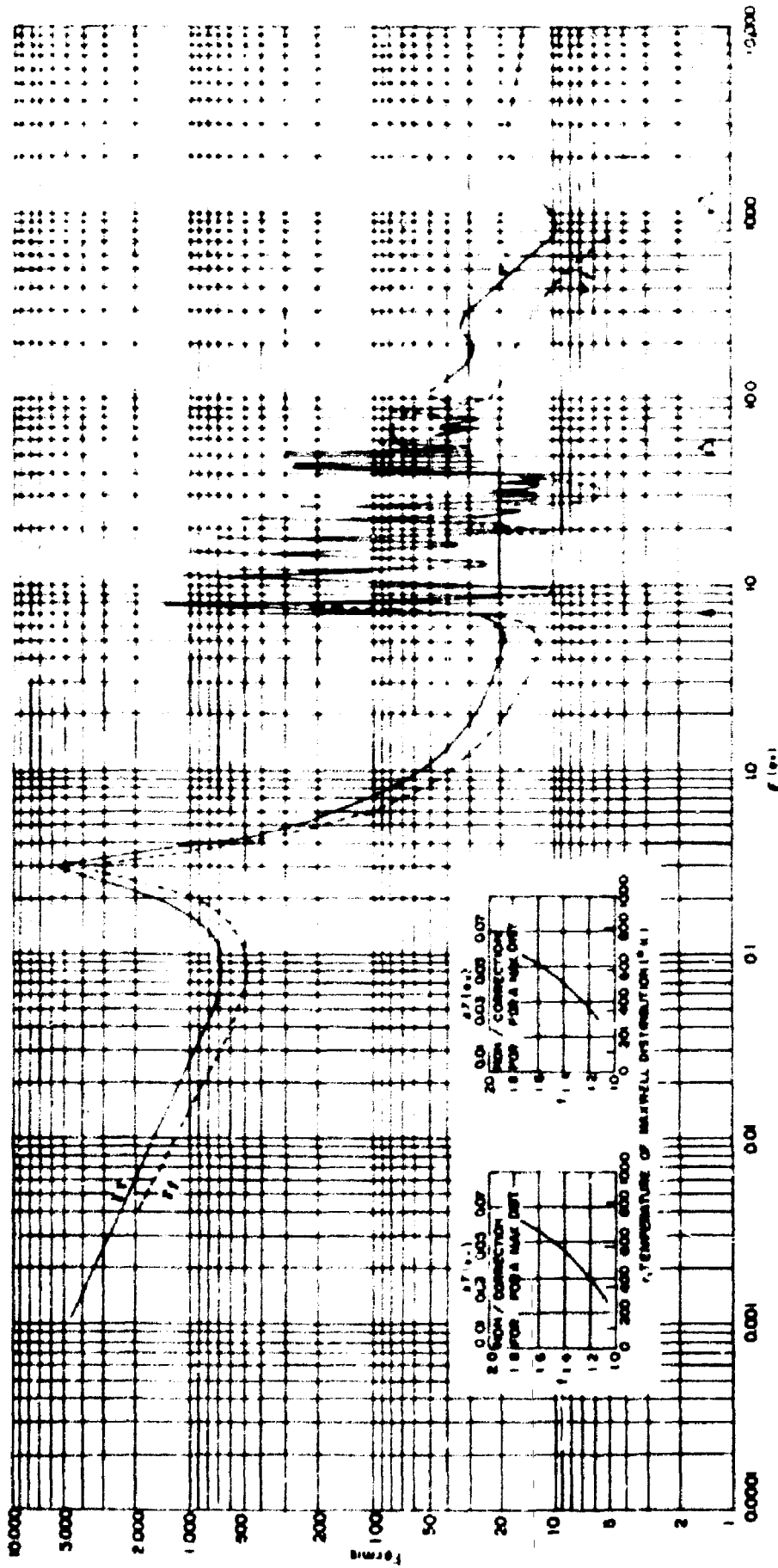


Fig 2 — Total ( $F_T$ ) and fission ( $F_f$ ) cross-sections of  $Pu^{239}$  in the low resonance and thermal regions

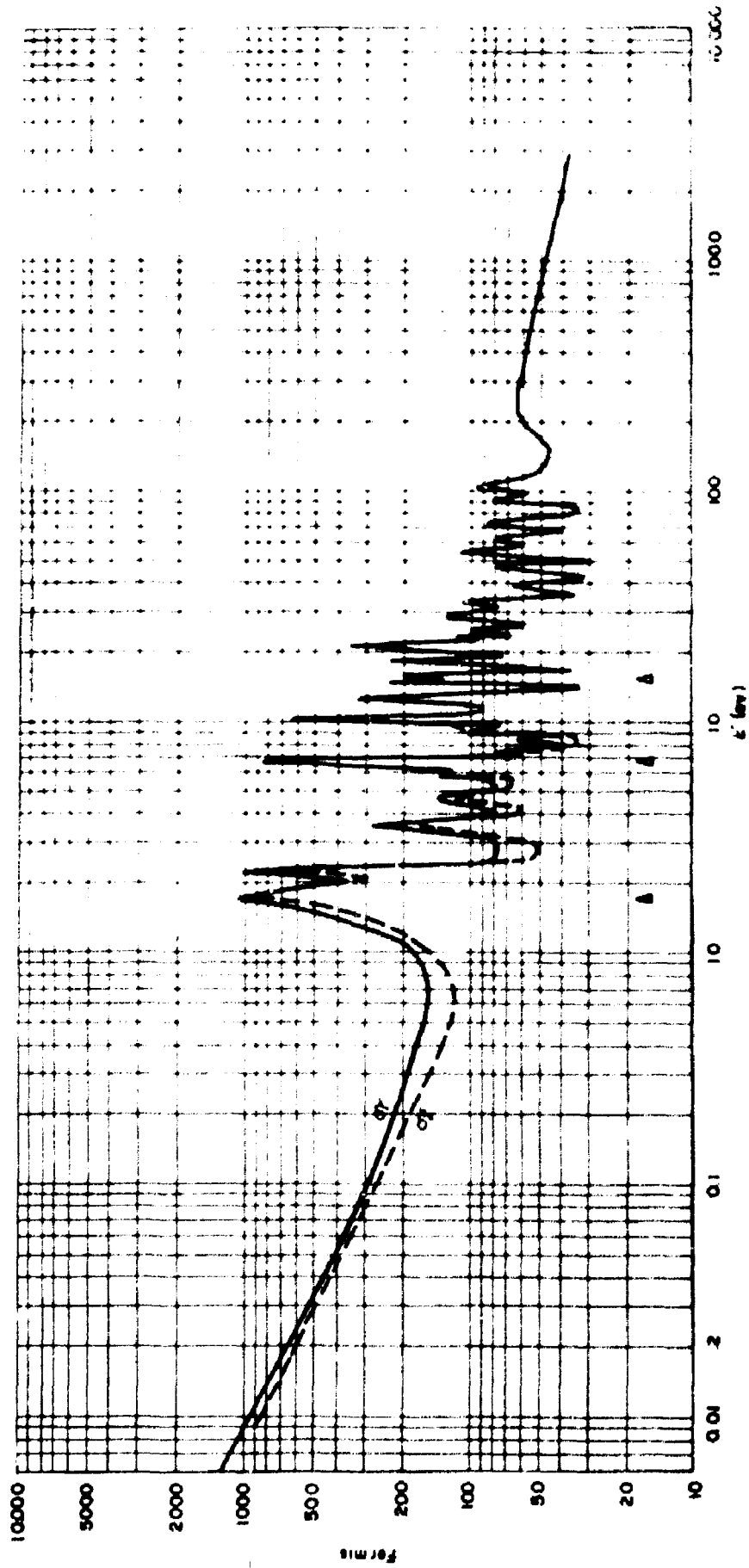


Fig. 3 - Total ( $\sigma_T$ ) and fission ( $\sigma_f$ ) cross-section of  $^{235}\text{U}$  in the low resonance and thermal regions

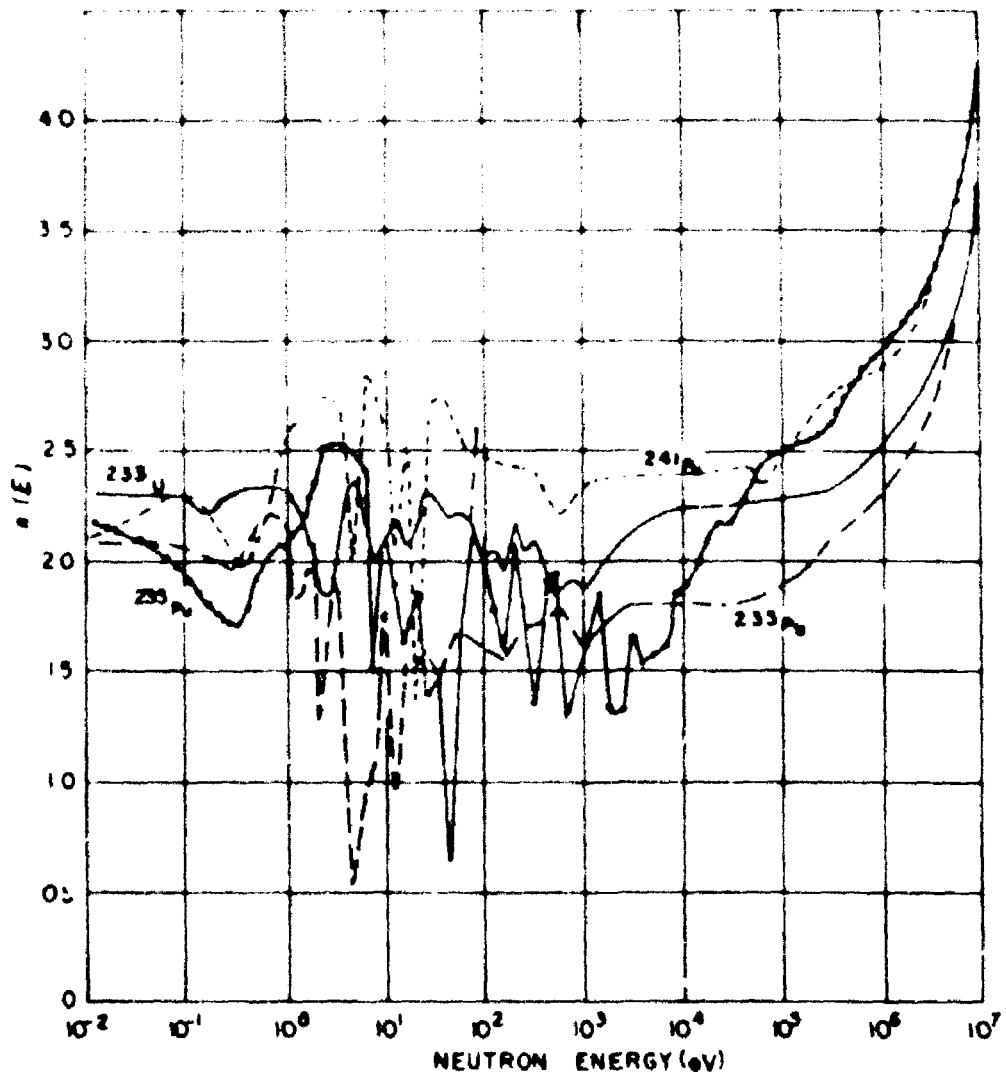


Figure 4. Energy dependence of eta for the principal fissile nuclides (4).

REPROCESSING OF THYR FULL ELEMENTS

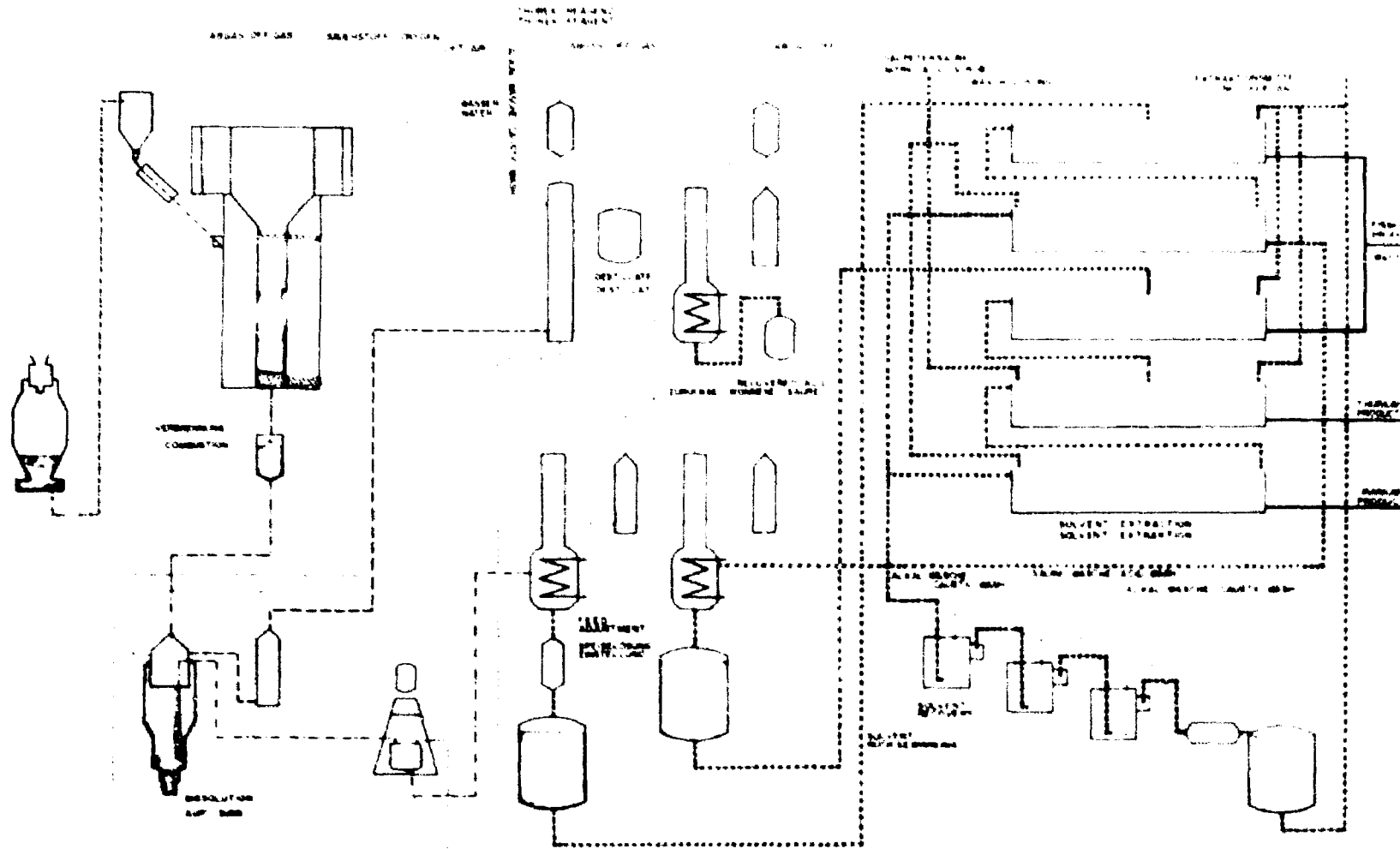
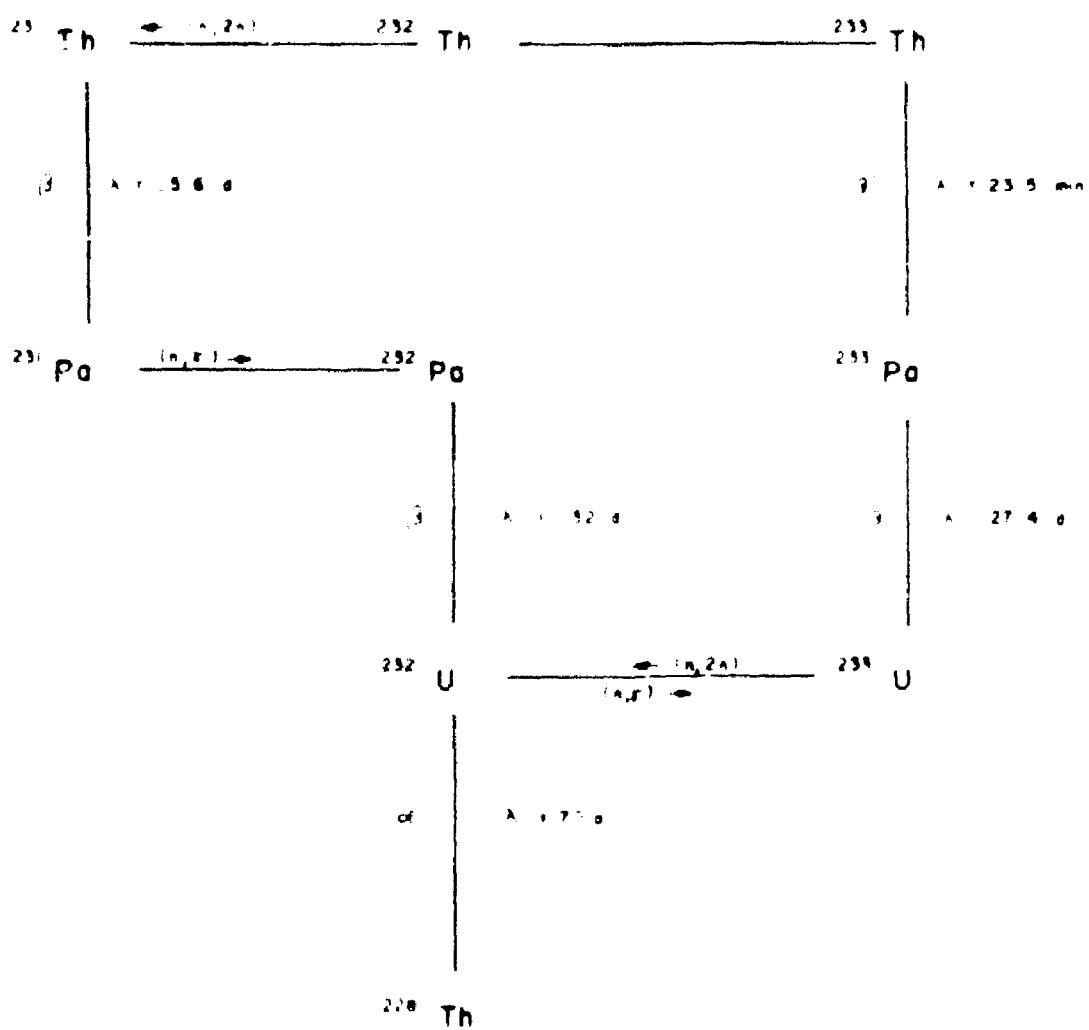


Fig 5 Generalized equipment flowchart of JUPITER

FIG 6 PRODUCTION OF U-232 and U-233



## Fabrication of Thorium Fuel elements

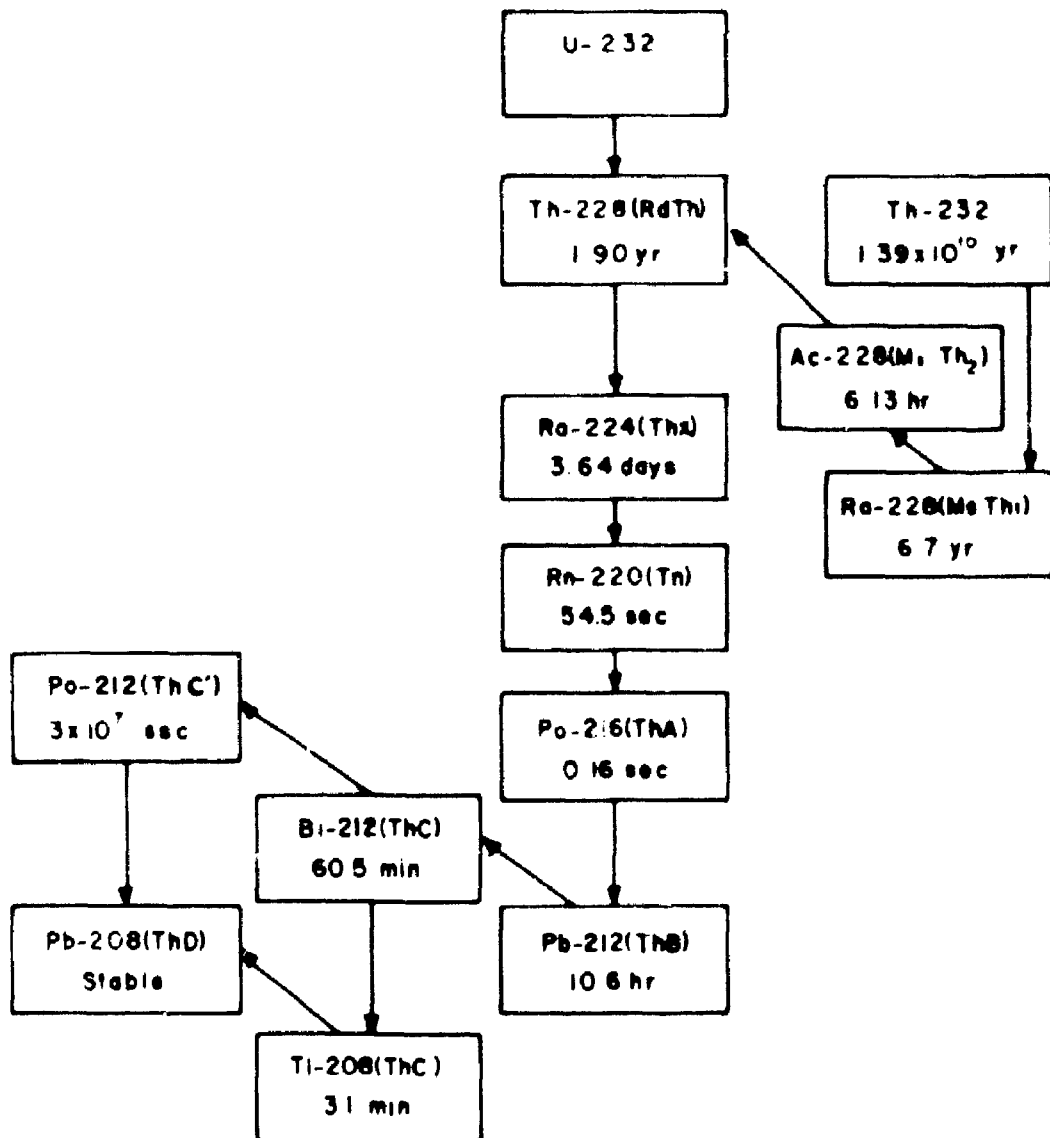


Fig 7. Decay scheme for thorium and U-232

Fabrication of Thorium Fuel Elements And

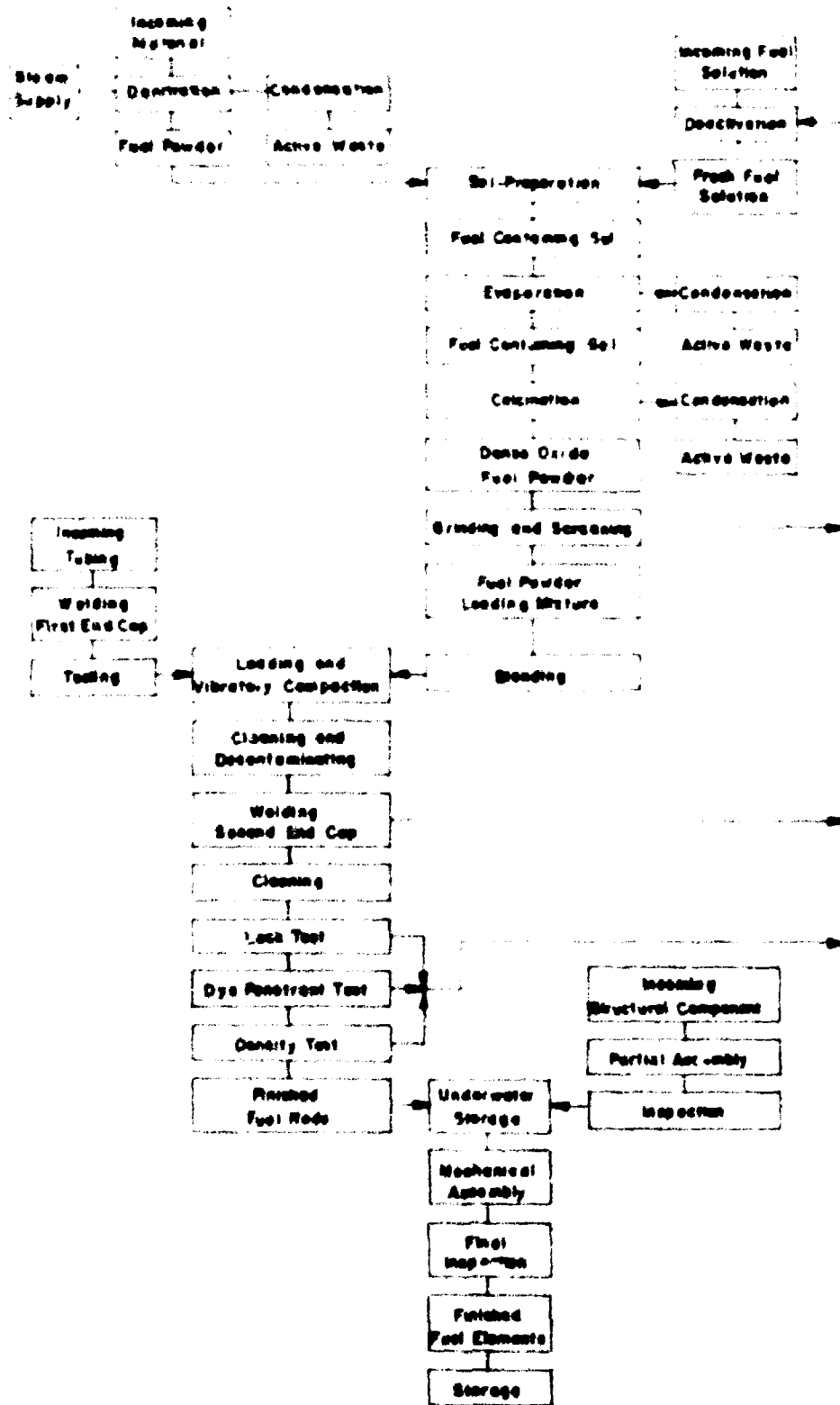


Fig. 6 Fuel fabrication flow sheet.



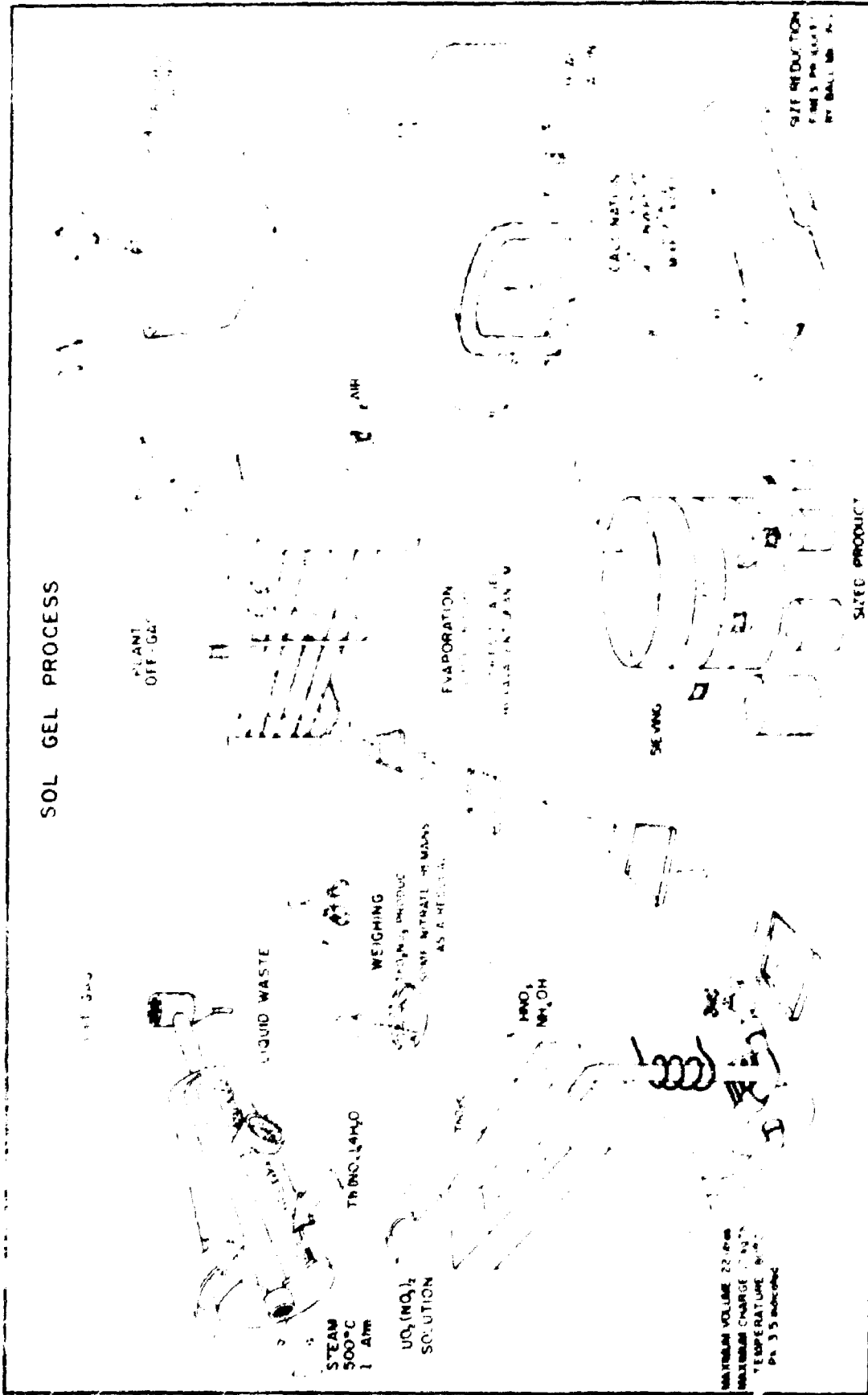


FIG 9 SOL GEL PROCESS FLOW SHEET

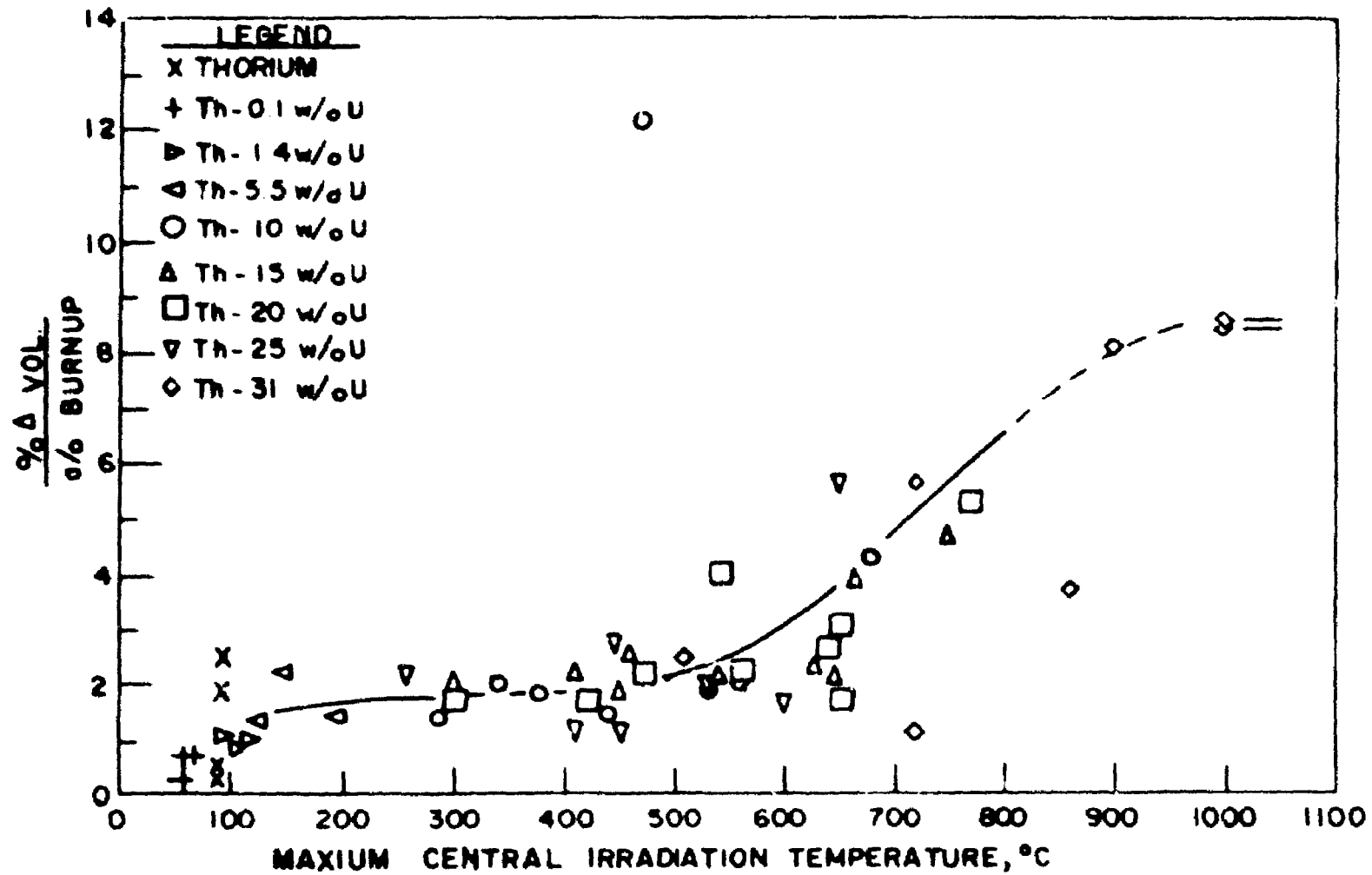


FIG. 10. Effect of irradiation temperature on the swelling rate of thorium and thorium-uranium alloys  
 Ref: KITTEL, J. H. et al., Effects of irradiation on thorium and thorium alloys. USAEC Rep. ANL 5  
 (1 April 1963)

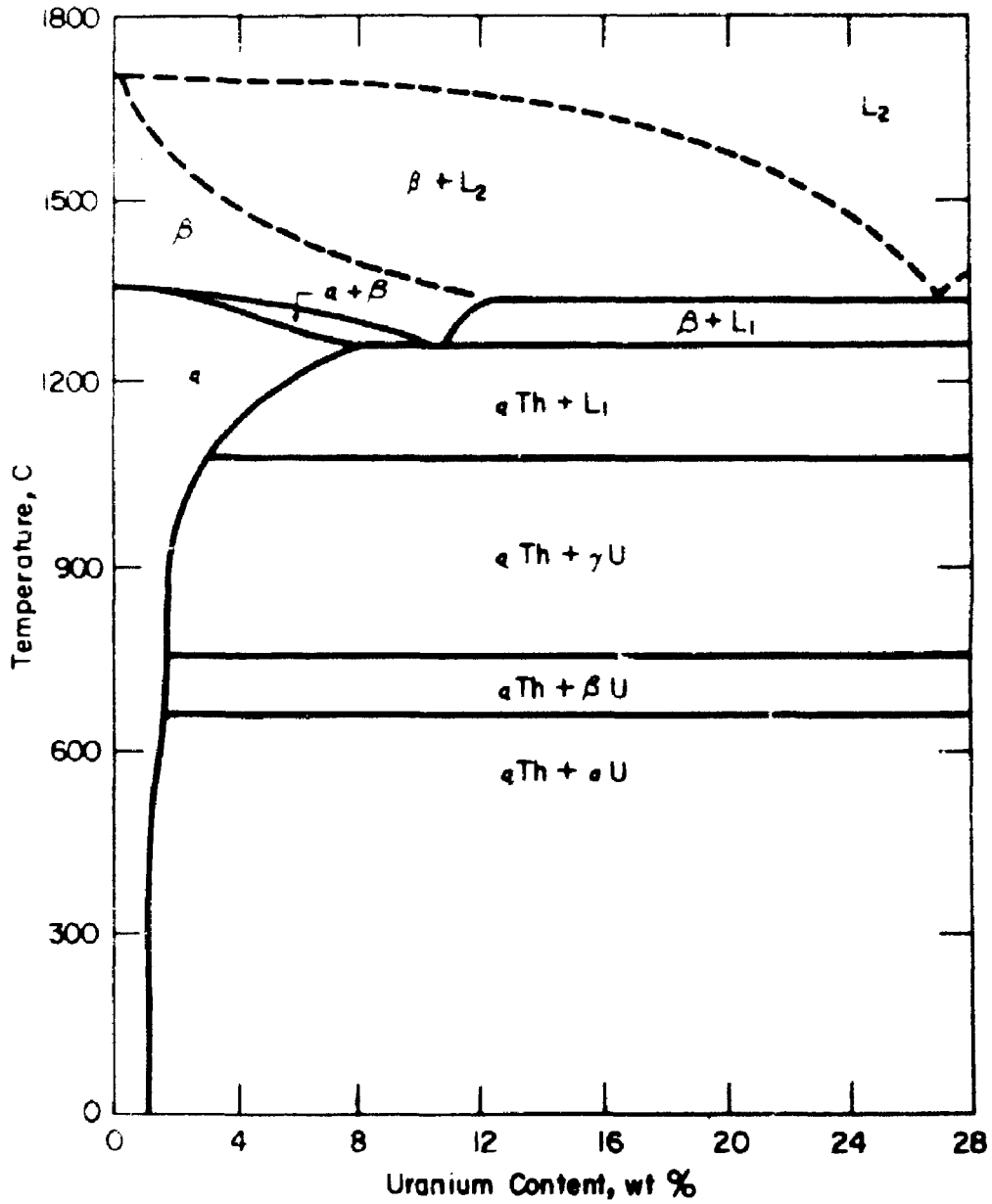


Fig. 11. Thorium-uranium equilibrium phase diagram for the range of interest of nuclear fuels.

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