

'NTERMOLECULAR VIBRATIONS IN THE LOW-FREQUENCY RAMAN SPECTRA OF L-ALANINE CRYSTAL

EUGENE LOH



INSTITUTO DE ENERGIA ATÔMICA Caixa Postal 11049 (Pinheiros) CIDADE UNIVERSITARIA "ARMANDO DE SALLES OLIVEIRA" SÃO PAULO — BRASIL

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Eugene Loh

Coordenadoria de Ciência e Tecnologia de Materiais Instituto de Energia Atômica São Paulo - Brasil

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ABSTRACT

The low frequency ≤ 150 cm⁻¹ polarized Reman spectra of L alanine single crystals by Weng and Storms can be interpreted from meaningfully in terms, instead of tranverse and longitudinal phonons by Wang and Storms of the intermolecular vibrations is ellibrations and/or translations of the four oriented alanine molecules in the unit cell of the orthorhombic crystal.

INTRODUCTION

The low frequency $\leq 150 \text{ cm}^{-1}$ polarized Raman spectra of L alanine single crystals have been measured by Wang and Storms¹ and interpreted¹ in terms of transverse and longitudinal phonons with, however, large orientational anistropy as well as polarization anomalies. To account for these unusual effects, a picture of dynamic disorder due to the protonic motion in the intermolecular hydrogen bonds between the NH³ group of one molecule and the CO² group of a neighboring molecule has been proposed¹.

In this note, we shall, by comparing the Raman intensities in the various polarized spectra, show that the polarized low frequency Raman spectra¹ of L alanine single crystals can be interpreted physically more meaningfully in terms, instead of transverse and longitudinal phonons¹ of the intermolecular vibrations² ³, i.e. librations and/or translations, of the four oriented alanine molecules within the unit cell of the orthorhombic alanine crystal, i.e. the orianted gas model.

ORIENTED GAS MODEL

To apply the oriented gas model, we need to approximate the three principle axes of the alanine molecule in the orthorhombic⁴⁻⁵ unit cell of alanine crystal by:

(i) u = 1 The projection of the u-axis on the ac crystallographic^{4,5} plane is taken to be the same as that from the C_{α} : C bond, where C_{α} is the α carbon atom and C is the carbon from CO_2^{-} . u is perpendicular to the subsequently defined v. The direction of u defined in this manner is close to both the crystallographic claxis^{1,5}, as indicated by the directional cosines in Table I, and the direction of the molecular chain^{4,5} in the crystal, u is, therefore, likely the "high" frequency axes, among u, v and w, of the intermolecular librations and translations

(ii) v - The direction of v is taken to be parallel to the line joining N of NH₃ and C_B of C_BH₃ in the alanine molecule and is close to that of the crystallographic⁴⁻⁵ a axis, Table 1.

(iii) w -- The direction of w is perpendicular to both u and v and is close to that of the crystallographic^{4,5} axis b, Table I

	8	b	C
u	0 322	- 0 202	0.923
v	0 935	0 225	0 275
w	- 0 152	0 952	0.262

Table I

Direction cosines of the approximated molecular axes u, v and w of the alanine molecule with respect to the crystallographic axes a, b and c of the orthorhombic L-alanine crystal.

Following the steps by Suzuki et al⁶ in deriving the relative Ramari intensities in various polarized spectra under the assumption of the oriented gas model:

The transformation matrices of molecules 1,2,3 and 4, in the $P2_12_12_1$ unit cell, between the molecule fixed co-ordinates u, v, w, and the crystal fixed coordinates a, b, c are expressed by

$$T_{1} = \begin{bmatrix} a_{u} & a_{v} & a_{w} \\ b_{u} & b_{v} & b_{w} \\ c_{u} & c_{v} & c_{w} \end{bmatrix}, \qquad T_{2} = \begin{bmatrix} a_{u} & a_{v} & a_{w} \\ -b_{u} & -b_{v} & -b_{w} \\ -c_{u} & -c_{v} & -c_{w} \end{bmatrix}$$
$$T_{3} = \begin{bmatrix} -a_{u} & -a_{v} & -a_{w} \\ b_{u} & b_{v} & b_{w} \\ -c_{v} & -c_{v} & -c_{w} \end{bmatrix} \text{ and } T_{4} = \begin{bmatrix} -a_{u} & -a_{v} & -a_{w} \\ -b_{u} & -b_{v} & -b_{w} \\ -b_{u} & -b_{v} & -b_{w} \\ c_{u} & c_{v} & c_{y} \end{bmatrix}$$
(1)

where a_{μ} is the cosine of the angle between the a and u axes and so on. The values of a_{μ} , a_{ν} ... are listed in Table I. The molecular tensor with respect to the crystal fixed coordinates is therefore given by

$$\alpha_{m_1} = T_{\mu} \alpha'_m \widetilde{T}_{\mu}$$
(2)

where i = 1, 2, 3 and 4 for molecules 1, 2, 3 and 4, repectively, and α'_m = derived pularizability tensor of the free molecule with respect to a normal vibration Q.

On the assumption of the oriented gas model we assume additive polarizabilities.

$$\alpha_{\rm u} = \alpha'_{\rm m1} Q_1 + \alpha'_{\rm m2} Q_2 + \alpha'_{\rm m3} Q_3 + \alpha'_{\rm m4} Q_4 \tag{3}$$

where α_u is the polarizability tansor of the unit cell. Using symmetry coordinates⁶ S_{Ag} , S_{Bg1} , S_{Bg2} and S_{Bg3} appropriate to the four-molecule unit cell of D_2 , we find

$$Q_{1} = (S_{Ag} + S_{Bg_{1}} + S_{Bg_{2}} + S_{Bg_{3}}) / \sqrt{4}$$

$$Q_{2} = (S_{Ag} + S_{Bg_{1}} - S_{Bg_{2}} - S_{Bg_{3}}) / \sqrt{4}$$

$$Q_{3} = (S_{Ag} - S_{Bg_{1}} + S_{Bg_{2}} - S_{Bg_{3}}) / \sqrt{4}$$

$$Q_{4} = (S_{Ag} - S_{Bg_{1}} - S_{Bg_{2}} + S_{Bg_{3}}) / \sqrt{4}$$

$$(4)$$

The derived polarizability tensors with respect to the unit cell vibrations, which can be obtained by taking the derivatives of equation(3) with substitutions of equations(1),(2) and(4;, are^7 :

$$\frac{\partial \alpha_{u}}{\partial S_{Ag}} = \sqrt{4} \begin{bmatrix} A_{ee} & 0 & 0 \\ 0 & A_{bb} & 0 \\ 0 & 0 & A_{cc} \end{bmatrix}, \qquad \frac{\partial \alpha_{u}}{\partial S_{Bg_{1}}} = \sqrt{4} \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & A_{bc} \\ 0 & A_{bc} & 0 \end{bmatrix}$$

$$\frac{\partial \alpha_{u}}{\partial S_{Bg_{2}}} = \sqrt{4} \begin{bmatrix} 0 & 0 & A_{ec} \\ 0 & 0 & A_{ec} \\ 0 & 0 & 0 \\ A_{ec} & 0 & 0 \end{bmatrix}, \qquad \frac{\partial \alpha_{u}}{\partial S_{Bg_{3}}} = \sqrt{4} \begin{bmatrix} 0 & A_{ab} & 0 \\ A_{ab} & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$
(5)

where the A's are functions of the direction cosines and the elements of the derived polarizability tensor of the free molecule.

Following Adamowicz et al⁸, we approximate the L-slanine molecule, which is of C_1 symmetry, by a C_2 molecule with the two fold axis almost parallel to u. The character table and selection rules for C_2 molecule indicate that the non-vanishing elements of the derived polarizability tensor of the free molecule are α'_{uu} , α'_{vv} , α'_{ww} , and α'_{vw} for the libration R(u) or the translational vibration T(u) about the u-axis. α'_{uw} and α'_{uv} for that about v and w axes. Thus, we obtain simple analytical expressions of A's for the vibrations of different species of the molecule. They are given in Table !!.

The squared values of A's, which are proportional to the expected intensities of a given laman peak in the various polarized spectra, are listed in Table III. The values in the columm of R(u), T(u) are more uncertain than that in the other two columns because of the presence of the α'_{uur} , α'_{vur} and α'_{uur} , in the former as shown in Table II.

Table III shows that for intermolecular vibrations about the u-axis, the orders of intensities are

$$(A_{1b})^2 > (A_{aa})^2 > (A_{cc})^2$$
 (10a)

4

and

$$(A_{ab})^2 > (A_{ac})^2 > (A_{bc})^2$$
 (10b)

and are satisfied by the "high" frequency peaks³, in Fig.1 and 2, at 105 cm⁻¹ and 113 cm⁻¹ for the meguality (10a) and at 105 cm⁻¹, 113 cm⁻¹ and 138 cm⁻¹ for inequality (10b)

For intermolecular vibrations about the v and w axes the intensity orders in the last column of Table III,

$$(A_{aa})^2 > (A_{cc})^2 > (A_{bb})^2$$
 (11a)

and

$$(A_{ac})^2 > (A_{bc})^2 > (A_{ab})^2$$
 (11b)

are satisfied by the "medium" frequency peaks¹ at 75 cm⁻¹ and 85 cm⁻¹ in Fig.1 and 2, for the inequality (11b), This indicates that these peaks are mostly of Bg mode in the D₂ factor group. However, for the lowered symmetry, e.g. C₂ factor group due to dynamic disorder¹, then these two "medium" frequency peaks¹ are also of Ag mode due to their strong intensities in (ac) spectra too.

The intensity orders in the next to last column are:

$$(A_{cc})^2 > (A_{bb})^2 > (A_{aa})^2$$
 (12a)

and

$$(A_{bc})^2 > (A_{ab})^2 > (A_{ac})^2$$
 (12b)

and are satisfied by the "low" frequency peaks at 48 cm⁻¹ and 40 cm⁻¹, in Fig 1 and 2, for both equalities (12a) and (12b). These two peaks are again the mixture of the Ag and Bg modes.

CONCLUSION

The oriented gas model seems to be able to explain the relative Ramen intensities in the polarized spectra of the low frequency vibrations in L alanine sigle crystals, which has hydrogen bonds as intermolecular bondings. The model assigns the intermolecular vibrations about the molecular axes u, v and w axes at (105 cm⁻¹, 113 cm⁻¹, 138 cm⁻¹), (75 cm⁻¹, 85 cm⁻¹) and (48 cm⁻¹, 40 cm⁻¹), respectively, of mixed A and B modes.

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	R(u), T(u)	R(v), T(v) 0	r R(w), T(w)
	$a_{ij}^{2}\alpha_{ijli}^{\prime} + a_{j}^{2}\alpha_{iji}^{\prime} + a_{ij}^{2}\alpha_{ijli}^{\prime} + a_{ij}^{2}\alpha_{ijli}^{\prime} + a_{ij}^{2}\alpha_{ijli}^{\prime} + 2a_{i}^{2}a_{ijli}\alpha_{ijli}^{\prime} + b_{ij}^{2}\alpha_{ijli}^{\prime} + 2b_{ij}b_{ill}\alpha_{ijli}^{\prime} + b_{ij}^{2}\alpha_{ijli}^{\prime} + 2b_{ij}b_{ill}\alpha_{ijli}^{\prime} + c_{ij}^{2}\alpha_{ijli}^{\prime} + c_{ij}^{2}\alpha_{ijli}^{\prime} + c_{ij}^{2}\alpha_{ijli}^{\prime} + 2c_{i}c_{ill}\alpha_{ijli}^{\prime}$	2ªuªwa′uw 2bubwa′uw 2cucwa′uw	2aua,au, 2bub,au, 2cub,au,
B P P C	$b_{u}c_{u}\alpha'_{uu} + b_{v}c_{v}\alpha'_{vv} + b_{w}c_{w}\alpha'_{ww} + (b_{v}c_{w} + b_{w}c_{v})\alpha'_{vw}$	(b _u c _w + b _w c _u)œ _{uw}	(b _u ς, + b _y c _u)α _ú ν
B ₉₂ ▲.c	aucuaiu + avcvaiv + awcwaiww + {avcw + awcv}oiw	(a _u c _w + a _w c _u)α _{úw}	(auc, + avcu)aúv
B ₉₃ A _{ab}	$a_{u}b_{u}\alpha_{uu}^{\prime} + a_{v}b_{v}\alpha_{vv}^{\prime} + a_{w}b_{w}\alpha_{ww}^{\prime} + (a_{v}b_{w} + a_{w}b_{v})\alpha_{vw}^{\prime}$	(aubw + awbu)œ́uw	(a,b, + a,b,)α,',

T

Table II

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Table III

Squared values of the elements of the crystal tensors

	R(u), T(u)	R(v), T(v) or	r R(w), T(w)
(A _{aa}) ²	$\sim 0.1 \alpha_{VW}^{\prime 2}$	$0.1 \alpha_{\rm uw}^{\prime 2}$	0.37 a'2
(A _{bb}) ²	>0 18 a'2	0.15 a'2	0.08 a'2
(A cc) ²	$\sim 0.02 \ \alpha_{vw}^{\prime 2}$	0.23 α ^{'2} _{uw}	0.26 α ^{'2} _{uv}
(A _{sc}) ²	$\sim 0.1 \alpha_{vw}^{\prime 2}$	0 03 α ^{′2} υw	0.6 α ^{′2} υν
(A _{ab})²	$\sim 0.74 \alpha_{VW}^{\prime 2}$	0.11 $\alpha_{uw}^{\prime 2}$	0.01 α ^{'2}
(A _{bc}) ²	$\sim 0.04 \alpha_{VW}^{\prime 2}$	0. 68 α ^{′2} _{UW}	0.07 a'2

I.



The Reman spectra of L-signine from 25 to 175 am⁻¹ for the (100) face and (010 - 'ece. (reproduced from reference 1)



Figure 2

The Raman spectra of L-slanine from 25 to 175 cm⁻¹ for the (001) face. (reproduced from reference 1)

RESUMO

Espectros Raman polarizados, de baixa frequência, < 150 cm⁻¹ de monocristais de Lialanina, obtidos por Wang e Storms podem ser reintérpretados, mais significativamento, em termos, ao invés dos fonons transverseis e longitudinais de Wang e Storms, das vibrações intermoleculares, i.e., "librations" e/ou translações, das quetro moléculas alarina orientadas na célula unitária do cristal ortorrômbico.

RESUME

Les spectres Remen polerisés, de bases fréquence, $\leq 150 \text{ cm}^{-1}$, de monocristaux de L alanine, obtenus per Wang et Storme peuvent être reinterpretés, plus significativement, en terms des vibrations intermioleculaires, cest-à-dire, des librations et/ou translations, des quatre molecules elanine orientées dens facellule unitaire du cristel orthorhombique.

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