

Directional correlations of γ transitions in ^{142}Ce

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Angular correlations of coincident gamma transitions have been measured in ^{142}Ce following the β^- decay of fission product ^{142}La ($T_{1/2}=93$ min) using a spectrometer consisting of a Ge and a Ge(Li) detector. The measurements were made on 14 direct and 13 skip cascades. Spin assignments were made to the levels at 2181 keV (3^+), 2727 keV (1^+), 3612 keV (2^+), and 4043 keV (2^+). In addition several previous spins assignments to other levels were confirmed. The multipole mixing ratios δ for 23 transitions were determined from the present results.

INTRODUCTION

Transitional nuclei with neutron numbers $84 \leq N \leq 90$ are of interest as they are expected to show gradual change from vibrational structure of a near spherical nuclei to the rotational behavior of the deformed region of rare-earth nuclei. The ^{142}Ce nucleus ($N=84$) thus provides a good opportunity for the study of the nuclear structure near magic number in this transitional region.

The β^- decay of ^{142}La to the levels in ^{142}Ce has been studied in the past by several workers.¹⁻⁵ The most complete study being that of Larsen *et al.*⁵ containing extensive Ge(Li)-Ge(Li) coincidence data. The levels in ^{142}Ce have also been studied by Coulomb excitation^{6,7} and $^{140}\text{Ce}(t,p)$ reaction.⁸

The measurements of γ - γ angular correlations in the decay of ^{142}La were performed by Prestwich and Kennett⁹ using two NaI(Tl) detectors and Basinger *et al.*¹⁰ using the combination of a Ge(Li) and six NaI(Tl) detectors. More recently Michelakakis *et al.*¹¹ and Wolf *et al.*¹² have used a combination of Ge(Li) detectors for these measurements. All these studies resulted in a level scheme of ^{142}Ce with spin and parity assignments made to several levels. The results are summarized in the Nuclear Data Sheets.¹³ Despite a large number of gamma cascades measured by Michelakakis *et al.*¹¹ several A_{kk} values are in serious disagreement with the results of Basinger *et al.*¹⁰ and Wolf *et al.*¹² leading to conflicting conclusions regarding the spin assignments to some of the levels and multipole mixing ratios of several γ transitions. In view of these discrepancies it was decided to investigate the angular correlations again using Ge(Li) and Ge detectors and at the same time to obtain data with improved counting statistics in order to better define some of the spin assignments to the levels made in the previous studies. The levels and transitions in ^{142}Ce were studied by measuring a total of 27 gamma cascades populated by the β^- decay of ^{142}La .

EXPERIMENTAL

The radioactive sources of ^{142}La were prepared by chemically separating the barium activity from the fission

products of uranium and later separating the lanthanum produced from the decay of barium. Approximately 1 g of uranyl nitrate hexahydrate were irradiated with thermal neutrons for 2 min at a flux of $\approx 5 \times 10^{12}$ n/cm²s in the IEA-R1 reactor at São Paulo. The barium activity was separated from the fission products a few minutes after the end of irradiation using a procedure similar to one described by Minkkinen.¹⁴ The barium chloride was finally dissolved in 5 ml of H₂O. The lanthanum carrier was added to this solution and lanthanum activity was allowed to grow for a period of 30 min, then separated from barium by precipitating La(OH)₃. The precipitate was centrifugally washed several times and finally dissolved in a drop of 1-M HCl. The dilute solution containing ^{142}La was transferred to a lucite source holder and taken to the gamma spectrometer for measurements. The source dimension was 2.5 mm \times 5 mm.

The γ - γ spectrometer consisted of a Ge detector with a volume of 89 cm³, maintained fixed, and a movable Ge(Li) detector with a volume of 45 cm³. The γ - γ coincidences were recorded using a standard low-noise fast-coincidence system and a 4096 channel pulse-height analyzer. The measurements were carried out at angles of 90°, 120°, 150°, and 180°. The angular position of the movable detector was changed every 20 min and the coincidence spectrum observed through the Ge detector was routed to an appropriate 1024 channel subgroup of the analyzer memory for each angle. Counting from a single source continued through a period of 160 min after which a fresh source with approximately the same activity was used. A total of 100 sources were used for the entire experiment.

The single-channel analyzer (SCA) window was set to accept the photopeak at 641 keV as seen in the Ge(Li) detector. An additional gate was placed adjacent to the main gate at the higher-energy side in order to determine the effects of Compton scattered radiation of higher-energy gamma rays included in the window setting. The intensities of coincident gamma rays were measured from the Ge-detector spectra at various angles and corrected for the source decay during the measurements—effects of Compton scattered radiation of higher energy and chance coincidences. The chance coincidences were determined

separately by introducing a delay of $1 \mu\text{s}$ in the signal pulses from one of the detectors before reaching the coincidence unit and recording the coincidence spectrum. The corrected photopeak areas were least-square fitted to the polynomial $W(\theta) = 1 + A_{22}P_2(\cos\theta) + A_{44}P_4(\cos\theta)$ to determine the angular correlation coefficients A_{kk} .

RESULTS

The direct γ -ray spectrum in the decay of ^{142}La obtained with the Ge detector is shown in Fig. 1(A). The γ - γ coincidence spectrum obtained with the 641-keV gate setting is shown in Fig. 1(b). The energy range is limited to 0–2.3 MeV and the coincidence spectrum represents the result of only a partial measurement not corrected for Compton contributions and accidentals. The angular correlation coefficients A_{kk} corrected for the finite solid-angle effects are given in Table I. The solid-angle corrections factors for the Ge detector were determined by numerical calculations¹⁵ and for the Ge(Li) detector they were taken from the tables of Camp and Van Lehn.¹⁶ The A_{kk} values for the gamma cascades measured by Basinger *et al.*,¹⁰ Michelakakis *et al.*,¹¹ and

Wolf *et al.*¹² are included in this table for comparison.

The multipole mixing ratios for the gamma transitions are presented in Table II. These values were calculated using spin sequence found most consistent with the observed angular correlation data, the decay properties⁵ and the results of reaction studies.^{6–8} The mixing ratios were determined by the usual χ^2 analysis as a function of δ for the mixed transition. The convention of Becker and Steffen¹⁷ was used for the definition of the mixing ratio. The multipole mixing ratios obtained in this study are compared with those of Basinger *et al.*,¹⁰ Michelakakis *et al.*,¹¹ and Wolf *et al.*¹² A partial level scheme of ^{142}Ce taken from the Nuclear Data Sheets¹⁹ is shown in Fig. 2. Only γ transitions of interest in this study are shown. The spin and parity assignments deduced from the present investigations are included in this figure. The parametric plots for some of the relevant spin sequences are shown in Figs. 3 and 4. The corrected values of A_{kk} coefficients with associated error bars have been displayed as (A_{22}, A_{44}) points in these plots. In the cases of skip cascades where the intermediate unobserved transition is the 895-keV transition a value of $\delta(895) = -0.63$

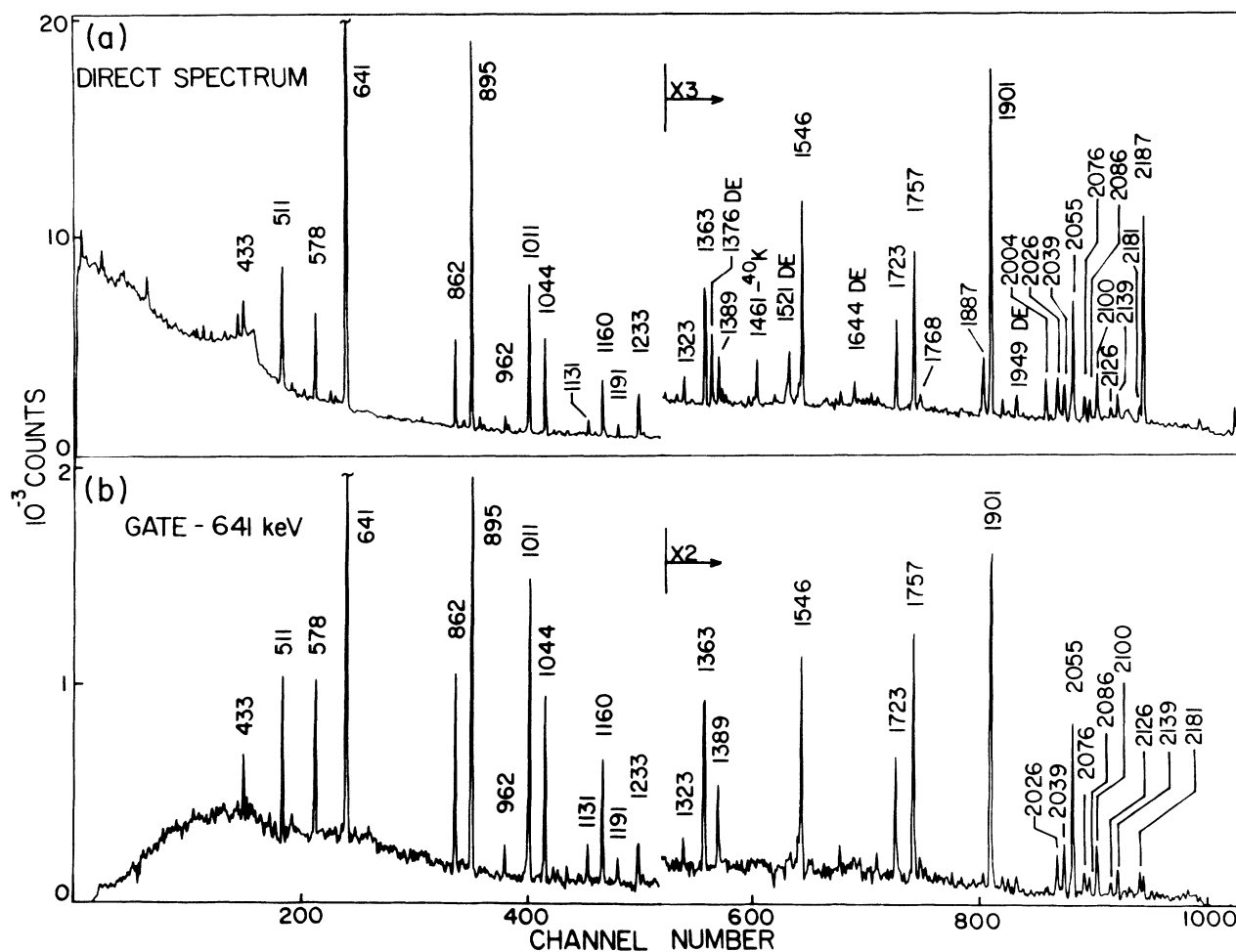


FIG. 1. (a) Direct γ -ray spectrum up to 2.3 MeV in the decay of ^{142}La observed with an HPGe detector and (b) γ -ray spectrum in coincidence with the photopeak 641 keV.

TABLE I. Results of directional correlation measurements of transitions in ^{142}Ce .

	Gamma cascade (keV)	A_{22}	A_{44}
1	578–641	0.102 ± 0.024	0.042 ± 0.037
		0.101 ± 0.021^a	-0.040 ± 0.030^a
		-0.055 ± 0.024^b	0.070 ± 0.043^b
		0.094 ± 0.055^c	-0.003 ± 0.060^c
2	895–641	0.510 ± 0.015	0.087 ± 0.025
		0.320 ± 0.020^a	0.010 ± 0.030^a
		0.325 ± 0.016^b	-0.020 ± 0.032^b
		0.417 ± 0.034^c	0.089 ± 0.036^c
3	1011–641	-0.074 ± 0.018	-0.006 ± 0.026
		-0.182 ± 0.022^b	0.006 ± 0.040^b
		-0.023 ± 0.045^c	-0.049 ± 0.050^c
4	1363–641	0.118 ± 0.027	-0.009 ± 0.042
		-0.053 ± 0.046^b	0.087 ± 0.084^b
		0.181 ± 0.044^c	-0.004 ± 0.047^c
5	1389–641	0.232 ± 0.105	0.868 ± 0.235
		0.182 ± 0.217^b	0.995 ± 0.447^b
		0.231 ± 0.281^c	1.436 ± 0.332^c
6	1546–641	-0.206 ± 0.024	-0.030 ± 0.034
		-0.295 ± 0.033^b	-0.008 ± 0.060^b
		-0.257 ± 0.045^c	-0.037 ± 0.049^c
7	1723–641	0.173 ± 0.035	-0.118 ± 0.054
		0.067 ± 0.039^b	-0.107 ± 0.070^b
		0.138 ± 0.055^c	-0.033 ± 0.062^c
8	1757–641	0.571 ± 0.025	-0.534 ± 0.041
		0.570 ± 0.070^a	-0.440 ± 0.110^a
		0.528 ± 0.056^b	-0.355 ± 0.097^b
		0.519 ± 0.043^c	-0.401 ± 0.047^c
9	1901–641	-0.176 ± 0.018	0.065 ± 0.026
		-0.130 ± 0.040^a	0.010 ± 0.050^a
		-0.223 ± 0.019^b	-0.031 ± 0.034^b
		-0.133 ± 0.024^c	-0.055 ± 0.026^c
10	2026–641	-0.353 ± 0.051	0.142 ± 0.071
		-0.343 ± 0.061^b	0.204 ± 0.110^b
		-0.291 ± 0.154^c	-0.056 ± 0.173^c
11	2055–641	0.512 ± 0.030	0.074 ± 0.046
		0.266 ± 0.039^b	0.057 ± 0.070^b
		0.455 ± 0.047^c	0.077 ± 0.053^c
12	2086–641	0.362 ± 0.079	-0.169 ± 0.123
		0.015 ± 0.128^b	-0.215 ± 0.276^b
13	2100–641	0.286 ± 0.052	-0.107 ± 0.081
		0.183 ± 0.053^b	-0.196 ± 0.100^b
		0.192 ± 0.085^c	-0.097 ± 0.108^c
14	2126–641	0.207 ± 0.196	0.841 ± 0.450
		-0.052 ± 0.254^b	1.022 ± 0.535^b
15	433–(578)–641	-0.216 ± 0.043	0.050 ± 0.062
16	962–(578)–641	0.261 ± 0.062	-0.089 ± 0.096
		0.152 ± 0.163^b	-0.016 ± 0.293^b
		0.346 ± 0.095^c	-0.094 ± 0.102^c
17	1363–(578)–641	0.168 ± 0.063	0.030 ± 0.098
		0.240 ± 0.060^a	0.085 ± 0.070^a
18	862–(895)–641	-0.082 ± 0.024	-0.014 ± 0.036
		-0.290 ± 0.041^b	-0.017 ± 0.074^b
		-0.029 ± 0.036^c	-0.004 ± 0.040^c
19	1131–(895)–641	0.007 ± 0.048	-0.162 ± 0.073
		-0.009 ± 0.116^b	-0.431 ± 0.202^b
20	1160–(895)–641	0.030 ± 0.027	0.013 ± 0.041
		0.041 ± 0.048^b	0.013 ± 0.087^b
		-0.016 ± 0.043^c	0.015 ± 0.049^c

TABLE I. (Continued).

	Gamma cascade (keV)	A_{22}	A_{44}
21	1191-(895)-641	0.056 ± 0.059 0.197 ± 0.191^b	0.034 ± 0.090 -0.195 ± 0.335^b
22	2076-(895)-641	0.174 ± 0.051 -0.090 ± 0.102^b	-0.063 ± 0.080 0.057 ± 0.192^b
23	2139-(895)-641	0.109 ± 0.066 -0.328 ± 0.141^b	0.094 ± 0.099 0.495 ± 0.275^b
24	2181-(895)-641	0.195 ± 0.074 -0.398 ± 0.141^b	0.091 ± 0.114 0.721 ± 0.287^b
25	1044-(1011)-641	-0.185 ± 0.021 -0.311 ± 0.031^b -0.143 ± 0.041^c	-0.022 ± 0.031 0.104 ± 0.056^b 0.057 ± 0.046^c
26	2039-(1363)-641	0.228 ± 0.050 -0.064 ± 0.078^b	-0.092 ± 0.077 0.034 ± 0.143^b
27	1233-(1546)-641	-0.138 ± 0.044 -0.252 ± 0.088^b	-0.007 ± 0.064 -0.030 ± 0.158^b

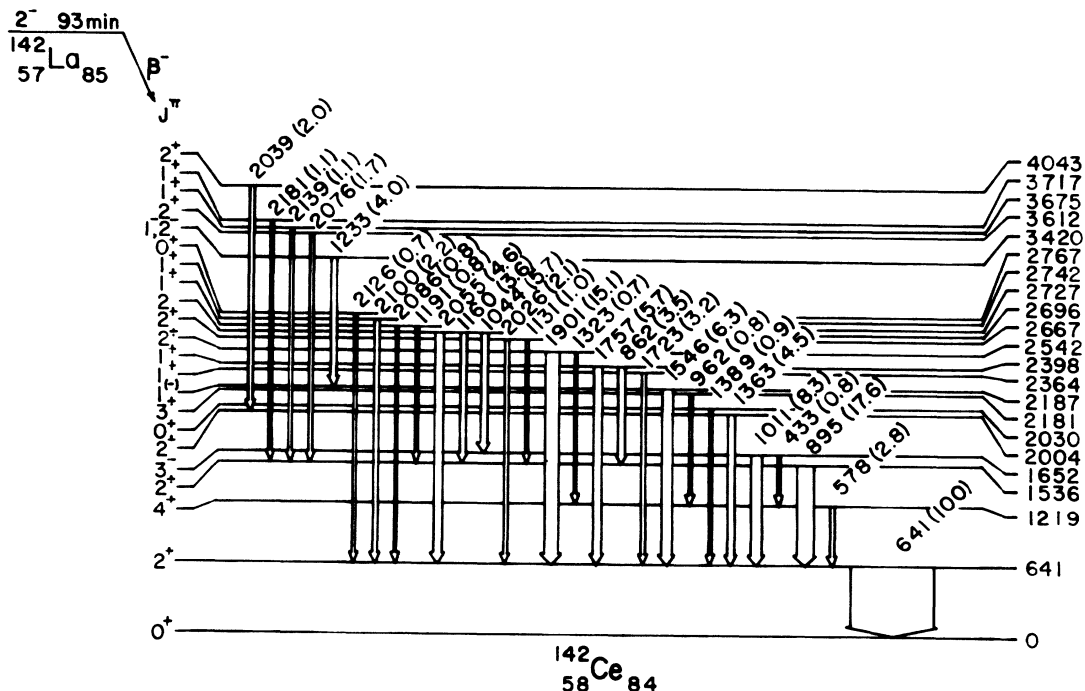
^aValues from Ref. 12.^bValues from Ref. 11.^cValues from Ref. 10.

has been used for the parametric plots.

The assignment of the ground state of ^{142}Ce is 0^+ as for all even-even nuclei and the 2^+ assignment for the 641-keV level follows from Coulomb excitation,^{6,7} the (t,p) reaction,⁸ and the decay studies.⁵ Since all gamma cascades studied involved the 641-keV γ ray as the lowest $2^+ \rightarrow 0^+$ transition, the spin assignment to most of the levels is relatively straightforward. The results of individual cascades and spin assignment to the levels are dis-

cussed below.

The spin and parity assignment of the 1219-keV level is believed to be $^{13}4^+$ which is consistent with most of the previous studies^{5,10,12} and also with the level systematics of the neighboring even-even nuclei. The only exception to this is the angular correlation result of Michelakakis *et al.*¹¹ for the 578-641-keV cascade which gives a 2^+ assignment. The present result for this cascade is in good agreement with the A_{kk} coefficients expected for the 4-2-

FIG. 2. A partial decay scheme of ^{142}La to the level in ^{142}Ce .

0 spin sequence.

The level at 1536 keV has a spin and parity¹³ of 2^+ . The present angular correlation result of the 895–641-keV cascade is in excellent agreement with this assignment.

The spin and the parity of the 1652-keV level is known to be 3^- .¹³ The results of the gamma cascades 1011–641 keV and 433–(578)–641 keV measured in the present work are in agreement with this assignment. The 1011-keV and 433-keV transitions are both predominantly dipole in character. The A_{22} coefficient for the 1011–641-

keV cascade obtained by Michelakakis *et al.*¹¹ although consistent with the 3-2-0 spin sequence is significantly different from the present value as well as that of Ref. 10.

The spin and the parity of the 2004-keV level is 2^+ .¹³ The present A_{kk} values for the 1363–641 keV-cascade are consistent with this assignment and agree with those of Basinger *et al.*¹⁰ but not with those of Michelakakis *et al.*¹¹

The 2030-keV level has been assigned a 0^+ spin and parity¹³ based on the angular correlation result of the 1389–641-keV cascade.^{10,11} As was pointed out earlier by

TABLE II. Multipole mixing ratios of gamma transitions in ¹⁴²Ce.

Level (keV)	Transition (keV)	$I_i^\pi - I_f^\pi$	Mixing ratio δ (This work)	Mixing ratio δ (Previous work)
1219	578	$4^+ - 2^+$	<i>E2</i>	
1536	895	$2^+ - 2^+$	-0.63 ± 0.10	$-0.11^{+0.02}_{-0.03}{}^b$ 0.61 ± 0.18^c
1652	1011	$3^- - 2^+$	-0.01 ± 0.03	-0.14 ± 0.03^b 0.06 ± 0.06^c
2004	433	$3^- - 4^+$	0.10 ± 0.06	
	1363	$2^+ - 2^+$	0.16 ± 0.04	0.41 ± 0.07^b -0.09 ± 0.06^c
2030	1389	$0^+ - 2^+$	<i>E2</i>	
2181	962	$3^+ - 4^+$	-0.56 ± 0.05	$1.01^{+2.08}_{-0.46}{}^c$
2187	1546	$1^{(-)} - 2^+$	-0.05 ± 0.05	0.04 ± 0.04^b $-0.01^{+0.04}_{-0.03}{}^c$
2364	1723	$1^+ - 2^+$	-0.38 ± 0.04	$-0.28^{+0.03}_{-0.04}{}^b$ 0.35 ± 0.05^c
2398	1757	$1^+ - 2^+$	-1.57 ± 0.10	$-0.93^{+0.28}_{-0.22}{}^b$ 1.06 ± 0.13^c
	862	$1^+ - 2^+$	0.03 ± 0.05	$\in (0.26, 0.36)^b$ 0.12 ± 0.12^c
2542	1901	$2^+ - 2^+$	0.65 ± 0.05	$0.55^{+0.40}_{-0.54}{}^a$ 0.02 ± 0.02^b 0.10 ± 0.30^c
	1323	$2^+ - 4^+$	<i>E2</i>	
2667	2026	$2^+ - 2^+$	1.33 ± 0.30	$\in (2.54, 1.02)^b$ -0.60 ± 0.05^c
2696	1131	$2^+ - 2^+$	$-5.56^{+1.99}_{-6.97}$	$\notin (-2.5, 2.95)^b$
	2055	$2^+ - 2^+$	-0.63 ± 0.10	-0.02 ± 0.06^b 0.55 ± 0.27^c
	1160	$2^+ - 2^+$	0.20 ± 0.05	0.22 ± 0.13^b -0.49 ± 0.30^c
	1044	$2^+ - 3^-$	0.02 ± 0.03	0.18 ± 0.04^b 0.03 ± 0.04^c
2727	2086	$1^+ - 2^+$	-0.60 ± 0.10	-0.23 ± 0.13^b
	1191	$1^+ - 2^+$	-0.43 ± 0.10	$\in (-0.23, -0.6)^b$
2742	2100	$1^+ - 2^+$	-0.50 ± 0.05	-0.04 ± 0.06^b 0.40 ± 0.09^c
2767	2126	$0^+ - 2^+$	<i>E2</i>	
3420	1233	$1^- - 1^-$	0.45 ± 0.05	$\in (0.48, 2.12)^b$
			or 2.24 ± 0.025	
		$2^- - 1^-$	0.40 ± 0.78 or $11.0^{+27.0}_{-5.0}$	$\in (0.44, 8.24)^b$
3612	2076	$2^+ - 2^+$	-0.68 ± 0.33	$\in (0.30, 1.31)^b$
3675	2139	$1^+ - 2^+$	-0.56 ± 0.10	$\in (0.89, 1.88)^b$
3717	2181	$1^+ - 2^+$	$-1.19^{+0.29}_{-0.45}$	$\in (0.65, 2.71)^b$
4043	2039	$2^+ - 2^+$	-0.99 ± 0.20	$\notin (0.28, -6.9)^b$

^aValues from Ref. 12.

^bValues from Ref. 11.

^cValues from Ref. 10 (the signs of mixing ratios have been changed to be consistent with the present sign convention).

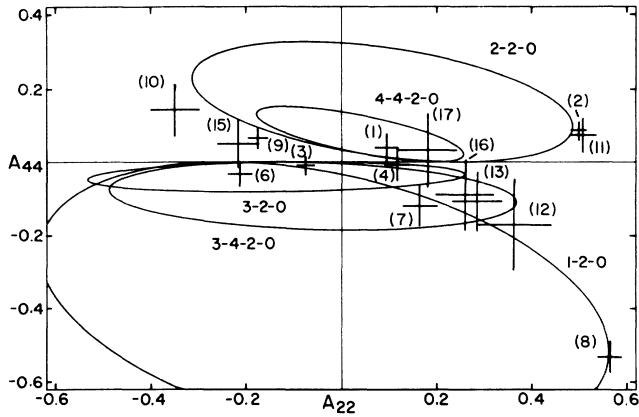


FIG. 3. Parametric plots for various spin sequences for direct cascades and skip cascades involving 578-keV intermediate transitions.

Basinger *et al.*,¹⁰ this cascade suffers an interference from the single escape of an intense 1901 keV γ ray and needs corrections. The present A_{kk} values are in fair agreement with both earlier results.^{10,11}

A spin and parity of 3^+ was indicated for the 2181-keV level from the angular correlation of the 962-(578)-641-keV cascade.¹⁰ Results of Michelakakis *et al.*¹¹ for this cascade allow 3^+ or 2^+ assignments despite the fact that authors assumed a spin and parity of 2^+ for the 1219-keV level in their analysis. The present result for the above cascade clearly indicates a 3-2-0 spin sequence. Since the 962-keV transition has a considerable quadrupole admixture the parity of the level is positive.

The spin of the 2187-keV level is 1 and its parity is probably negative.¹³ The present result for the 1546-641-keV cascade is quite consistent with the 1-2-0 spin sequence. Since the 1546-keV transition is almost pure dipole, nothing can be said about the parity of the

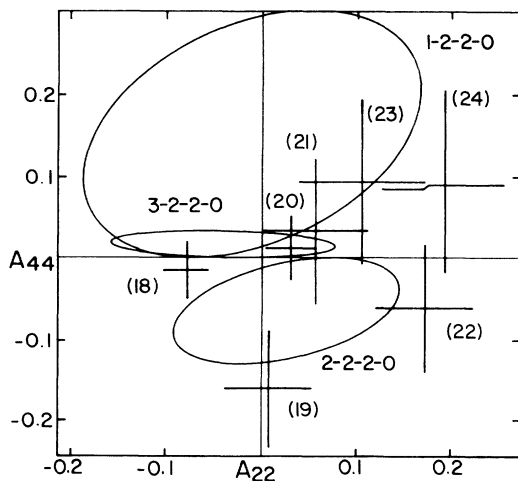


FIG. 4. Parametric plots for various spin sequences for the skip cascade involving 895-keV intermediate transitions.

level.

The spin of the 2364-keV level is indicated as 1 or 3 from the present result as well as the previous result¹¹ of the 1723-641-keV cascade. However the presence of a crossover transition to the ground state limits the spin value to 1. Since the 1723-keV transition shows a substantial quadrupole admixture the parity of the level is positive.

The spin and the parity of the 2398 keV level is¹³ 1^+ based on the previous angular correlation results of the 1757-641-keV cascade.¹⁰⁻¹² The present A_{kk} values for the cascade are in excellent agreement with the earlier results¹⁰⁻¹² and confirm the 1^+ spin assignment. The present result for the 862-(895)-641-keV cascades lends an additional support for this assignment.

While Larsen *et al.*⁵ assigned a spin and parity of 2^+ for the 2542 level from the decay study, the results of the (t,p) reaction⁸ and angular correlation of the 1901-641-keV cascade^{10,11} suggested a spin of 1 for this level. As discussed earlier the spin of the 1219-keV level is now known to be 4^+ and since the 2542-keV level decays to the 1219-keV level through the 1323-keV transition, an assignment of spin 1 for the 2542-keV level is ruled out. The present results for the 1901-641-keV and 1323-(578)-641-keV cascades are in good agreement with those of Wolf *et al.*¹² and clearly support a 2^+ assignment for the 2542-keV level.

A spin and parity of 2^+ is known for the 2667-keV level.¹³ The present angular correlation results of the 2026-641 keV and 1131-(895)-641-keV cascades are in agreement with the previous results^{10,11} and clearly indicate a spin of 2^+ for the level.

The angular correlation results of three gamma cascades: 2055-641 keV, 1160-(895)-641 keV, and 1044-(1011)-641 keV serve to define the spin of the 2696-keV level. The result of the 2055-641-keV cascade are quite unambiguous indicating the 2-2-0 spin sequence. However contrary to the results of Michelakakis *et al.*¹¹ the quadrupole admixture of the 2055-keV transition is much higher suggesting a positive parity for the level. The results of 1160-(895)-641-keV and 1044-(1011)-641-keV cascades are also consistent with the above spin assignment.

Two gamma cascades 2086-641 keV and 1191-(895)-641 keV deexciting the 2727-keV level were measured in the present study. The result of the 2086-641-keV cascade provides a clear indication of the 1-2-0 spin sequence. Since the 2086-keV transition has a considerable quadrupole admixture the parity of the level is positive. The result of the 1191-(895)-641-keV cascade although consistent with this assignment cannot exclude other spin values.

The spin of the 2742-keV level is 1,¹³ derived from the angular correlation results of the 2100-641-keV cascade.^{10,11} The present results for this cascade are in very good agreement with this assignment. The multipole mixing ratio of the 2100-keV transition $\delta(2100) = -0.50 \pm 0.05$ shows at least 17% quadrupole admixture suggesting a positive parity for the level.

The spin and parity of the 2767-keV level is known to be 0^+ (Ref. 13) based on the angular correlation of the

2126–641-keV cascade.¹¹ The present A_{44} value for this cascade although it does not cover the theoretical value expected for the 0-2-0 sequence, is large enough to exclude other spin sequences.

Michelakakis *et al.*¹¹ suggested the spin and parity of the 3420-keV level as 0^- , 1^- , or 2^- from their angular correlation result of the 1233–(1546)–641-keV cascade. The present result for this cascade is consistent with 1^- or 2^- spin and parity assignments but not 0^- . Mixing ratios of the 1233-keV transition for the respective spin assignments are given in Table II.

The spin and parity of the 3612-keV level was limited to 1^\pm or 2^+ from the angular correlation of the 2076–(895)–641-keV cascade.¹¹ The present A_{kk} values for this cascade differ considerably from those of Michelakakis *et al.*¹¹ (Table I) and clearly indicate a spin of 2 for the 3612-keV level. Substantial quadrupole admixture of the 2076-keV transition ($\delta = -0.68 \pm 0.33$) suggests a positive parity for the level.

The spins and the parities of the 3675- and 3717-keV levels are known to be 1^+ .¹³ The present A_{kk} values for the 2139–(895)–641- and 2181–(895)–641-keV cascades although quite different from those of Michelakakis *et al.*¹¹ clearly indicate a 1-2-2-0 spin sequence. The present results therefore confirm 1^+ assignment for both levels.

The angular correlation of the 2039–(1363)–641-keV cascade measured by Michelakakis *et al.*¹¹ allows any of the 1^\pm , 2^\pm , 3^\pm possible values for the spin and parity of the 4043-keV level. The present results for the above cascade are quite unambiguous, indicating a spin of 2 for the level. Since the quadrupole admixture of the 2039-keV transition is quite large the parity of the level is positive.

DISCUSSION

The low-lying excited states of even-even nuclei in the mass range $60 \leq A \leq 150$ are usually described as vibrations about a spherical equilibrium shape, as rotations of a soft deformed core or as excitations of two particles (quasiparticles) from ground state. These various models lead to quite different predictions for the electromagnetic transition probabilities. A systematic investigation of quantities such as the multipole mixing ratios $\delta(E2/M1)$ for γ transitions is therefore important in providing an understanding of the structure of these nuclei in terms of collective or single-particle effects.

Considering that the nucleus of ^{142}Ce has only two neutrons more than the closed shell of $N=82$ one might attempt to explain some of the energy levels in terms of a vibrational model. In this context the 641-keV (2^+) level would be the one-phonon state and levels at 1219-keV (4^+) and 1536-keV (2^+) would be the members of a two-phonon triplet (0^+ , 2^+ , 4^+). The 0^+ level at 2030 keV is however too high in energy to be the member of two-phonon triplet, being rather more appropriate for a three-phonon state. It may be noted that the 1536-keV (2^+) level decays exclusively to the 641-keV level, with no crossover transition. The absence of a crossover transition to the 0^+ ground state is surprising. This transi-

tion, although forbidden in the simple vibrational model, usually occurs with reduced intensity (at least a few percent of that of the stopover transition). The mixing ratio of the 895-keV- $2_2^+ \rightarrow 2_1^+$ transition was deduced to be -0.63 ± 0.10 . The corresponding transitions in the neighboring $N=84$ isotones ^{140}Ba ($\delta = -1.1 \pm 0.14$) (Ref. 18) and ^{144}Nd ($\delta = -1.6 \pm 0.5$) (Ref. 19) show somewhat larger quadrupole content but have the same sign for the mixing ratio.

The 3^- state at 1652 keV is known to be of a collective nature⁶ and most probably corresponds to the one-phonon vibration. Identification of the three-quadrupole-phonon states (0^+ , 2^+ , 3^+ , 4^+ , and 6^+) is difficult. The levels at 2004 keV (2^+), 2030 keV (0^+) and 2181 keV (3^+) have the correct energies for the three-phonon states; however, the decay characteristics of the 2^+ and 0^+ states suggest that they probably do not correspond to the three-phonon multiplet.

The main difficulty in this type of interpretation arises when one considers the electromagnetic properties of the γ transitions between the levels. The present results show that a large number of transitions in ^{142}Ce have considerable $M1$ admixture and these are difficult to explain in terms of simple vibration model. In particular the $2_2^+ \rightarrow 2_1^+$ transition is forbidden as $M1$ if the states are considered pure phonon states. It appears, therefore, that the interpretation of the excited states of ^{142}Ce in terms of a simple vibrational model is of limited value.

The dynamic deformation theory based on the pairing-plus-quadrupole model has been employed recently²⁰ for a detailed study of the $N=84$ vibrational nucleus ^{144}Nd . The multipole mixing ratios $\delta(E2/M1)$ have been compared with available data. The calculated $\delta(E2/M1)$ values for nearly all $2_n^+ \rightarrow 2_1^+$ transitions are, however, large indicating pure $E2$ in disagreement with the experimental results. The large $M1$ strength in ^{144}Nd was attributed by the author²⁰ to the possible two-quasiparticle admixture in several 2^+ states near 2 MeV not explicitly included in the calculations.

In recent years the neutron-proton version of the interacting boson model²¹ (IBM-2) which distinguishes between neutron and proton bosons has been quite successful in describing the properties of low-lying collective states in even- A nuclei of medium and heavy mass.^{22,23} Because of extra neutron-proton degrees of freedom IBM-2 can give rise to collective $M1$ transitions through the mixing of states with different neutron-proton symmetry. In a geometrical picture these mixed F -symmetry states correspond to slightly different deformations of neutron and proton cores or to the isovector vibrations.²⁴ Hamilton *et al.*²⁵ obtained reasonably good agreement between the calculated and experimental values of $\delta(E2/M1)$ for $2_3^+ \rightarrow 2_1^+$ transitions in $N=84$ isotones, ^{140}Ba , ^{142}Ce , and ^{144}Nd , where they identified the 2_3^+ states in these nuclei as the mixed F -symmetry states. However it was recognized²⁰ that a clear distinction between the mixed F -symmetric and two-quasi-particle states based only on $M1$ strength and energy value is rather difficult.

In the present work angular correlations of 27 gamma cascades were measured and multipole mixing ratios of

23 transitions were determined. Spin and parity assignments to the levels at 2181 keV (3^+), 2727 keV (1^+), 3612 keV (2^+), and 4043 keV (2^+) are suggested. In addition many previous spin assignments to other levels were confirmed. It is our hope that the results of the present investigation will be useful for some future calcu-

lation in an attempt to explain the level structure and electromagnetic transitions in ^{142}Ce .

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