HYDROGEN TRANSPORT PROCESS IN OH DOPED KC1 AND KBr*

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It was investigated the kinectics of formation and destruction of hydrogen related defects originated from the randon aggregation of hydrogen atoms in their ground state that were allowed to diffuse through the crystal after the photodissociation of OH⁻ defects. Formation and ac tivation energies for the several processes involved in three main steps were obtained utilizing 1st. order kinectics. A model for a new U_{2x} center is proposed.

In a preliminary study on the kinectics of U2to H2O defects conversion after the photo dissociation of OH ions in alkali halides it was shown¹ that the resulting hydrogen atom diffuses in the dark with a thermally assisted motion that brings this atom into two different configurations. In one configuration the hydro gen atom remains trapped by an OH ion forming a sort of H_1^0 - OH^- association (U_{2X}center) that produces electronic absorptions in the UV (45000 cm^{-1} for KC1 and 40000 cm^{-1} for KBr). The other configuration is the H20⁻ defect itself that is produced by a direct and definitive trapping of the H_1^0 atom by one OH ion forming with it a water molecule in a substitutional position and imbebed in an electronic cloud also with inde pendent electronic absorption^{2,3}. These two con figurations are produced via reaction channels that compete with each other at the expenses of available and mobile hydrogen atoms and substitu tional static OH⁻ impurities. In the sequence of creation and anihilation processes that $\stackrel{\circ}{}$ originated from the photodissociation of the OH- ion, be fore the nearly complete restoration of the ori ginal OHT concentration, F and U centers are also formed. In this work we investigate the me chanisms behing these atomic and ionic transport processes that take place with the participating centers in their ground state.

The experimental procedures followed stand ard techniques of optical absorption at low temperatures. Typical $10x5x1 \text{ mm}^3$ samples were cleaved from boules of KCl and KBr doped with KOH at 10^{-4} mole % grown in our laboratory by a Zochralski method from U.P. grade material. Op tical absorption measurements were performed with a Cary 17D spectrophotometer utilizing an opti cal cryostat from Janis, model 8DT provided with a gas exchange chamber as a thermal switch. Tem perature was controlled and measured with copper -constantan thermocouples referred to an electron ic ice point reference. The photodissociation of OH⁻ was done with a resonant wavelenght provided by a xenon lamp plus a monochromator. The samples were pulse annealed under a programmed routine to allow a reproducible dT/dt shape and then cooled to LNT for the optical measurement.

The thermally assisted formation and de struction processes were analysed utilizing first order kinectics where applicable, following the procedure done by Fritz⁴ and under the as sumption that after the initial dissociation of the OH⁻ ion, only one mobile entity ⁵ is responsible for the formation of one or more in dependent defects, after being trapped at different sites. In figures 1 and 2 there are shown for KC1 and KBr respectively the behaviour of the absorption bands maxima of the defects involved plotted against the annealing temperature and after the initial dissociation of the OH⁻ . In these sequences of destruction/ creation of defects that start with the destruc tion of the U2 center, several intermediate de fects appear and disappear. In both cases these sequences end up with F and U centers. After prolonged stay of the sample in the dark at room temperature, these two centers also disappear and restore the OH concentration to almost its initial value.

The full cycle of defects transformation was divided into three main steps, each of them characterized by the thermal extinction of one defect. These steps occured when the extinction of U_2 , U_{2x} and H_20^- took place. The temperatures of maximum chance (dn/dT = max) in dicated in table 1 correlate precisely the ex tinction of one defect with the appearance of one or more defects, at the expenses of the destroyed one. In this table there are shown activation energies obtained by fitting the three regions of interest from curves of figu res 1 and 2 with a first order kinectics. Below the activation energy values it is shown the percent population of centers destroyed or cre ated. These percentual values represent the probabilities that a certain reaction channel occurs, taking into consideration the concen tration and oscillator strengths of the centers

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Figure 1 - Absorption bands maxima at 77 K for pul se annealed KCL:OH⁺+ U₂ system after the photodissociation of OH⁻. In the final step U centers are preferentially formed.



Figure 2 - Absorption bands maxima at 77 K for pul se annealed KBr:OH⁻ + U₂ system after the photodissociation of OH⁻. In this case F centers are preferentially for med in the final step.

involved as done previously¹. It is important to note that these probability values are obtained independently from the activation energies that resulted from the fitting of the kinectic curves.

In the first step the extinction of the U2 center leads to the formation of U_{2x} and $H_2O^$ centers. The temperature dependence of the U2 concentration yielded activation energies of 0.10 eV for KCl and 0.13 eV for KBr. This pure ly atomic process involving thermal migration and trapping of the hydrogen atom, takes place in the dark and consequently with this atom in its ground state. In this way a U, center which is forced to move from its stable inters titial position by a phonon mechanism, produces U, related defects without lattice Jahn-Teller distortion or crowdion creation that are ob served by the photo-excitation of the U2 center at low temperatures⁶. Consequently distinct defects are produced that do not change the bas ic lattice defect configuration since no pola rization effects due to optical excitation change the lowest empty 4s type orbital (K⁺) of the conduction band. These low activation ener gy values for the hydrogen atom to escape (while in its ground electronic state) reflect its small size and charge neutrality and, conse quently, a shallow potential well where it was trapped. In spite of diffusing in its ground state the hydrogen atom has low energy values that resemble photon assisted process of ionic diffusion of point defects in their excited sta te, what is well known for lowering the step $di\overline{f}$ fusion activation energies of centers such as the F_A^{7} . the F_A

The formation energies necessary to produce U_{2x} and H_2O^- centers indicate in KBr that since U_{2x} centers have a lower forma tion energy than H₂O⁻ centers (0.15 vs 0.19 eV) U_{2X} centers are preferentially formed and in larger number (for KCl the U_{2x} is not a resolved band so no correlation was possible). This was consistently observed through the num ber of centers produced after the 100% destruc tion of U₂ that produces 34% of U_{2x} and 16% of H₂O⁻ centers. The remaining 50% of U₂ centers not accounted for are possibly aggregated into H₂ molecules that are invisible by optical techniques. In this step, the rates of formation involved are dependent on the available OHdefects left without being photodissociated to provide the necessary number of traps for U_2 centers to aggregate. In this work the initial OHT photodissociation was limited to 35% of the original OH concentration. It should be pointed out that a certain number of U2 centers react back with O⁻ defects producing OH⁻ centers. This back process was verified through a thermally assisted luminescence⁸ that is produced by the de-excitation of OHT defects formed in their excited state after the capture of the hydrogen atom by the substitutional 0⁻ defect left by the initial OHT photodissociation. By this way the number of available OH T traps is in fact larger than the quantity left after the OH pho todissociation. To the statistically available OH traps one should consider the dynamically produced ones that will increase the total number

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Table 1. Destruction and formation energies of the different centers involved in this study, for KCl and KBr in the three main temperature steps. The temperature where dn/dt is maxi mum is indicated in parenthesis for each crystal. The numbers in percent correspond to the amount of centers involved in each step. Energies are in eV.

Step	Crystal (T)	Destruction Energy			Formation Energy			
		U ₂	U _{2x}	H ₂ 0 ⁻	U _{2x}	H ₂ 0 ⁻	F	U
I	KC1 (100K)	0.10 100%			0.19 ?	0.19 ?		
	KBr (95K)	0.13 100%			0.15 34%	0.19 16%		
II	KC1 (140K)		0.18 100%			0.19 100%		
	KBr (205K)		0.36 100%			0.34 55%		0.37 45%
III	KC1 (230K)			0.55 100%			0.80 14%	0.51 57%
	KBr (245K)			0.46 100%			0.46 100%	

of traps for the hydrogen atom to interact with. What happens next to a U $_{2\,\rm X}$ centers is its final transformation into a $\rm H_2O^{-}$ center. For KCl this conversion is complete what is justified in part by a lower temperature of stabilization and a smaller activation energy when compared to KBr where the production of U centers also results as a competitive reaction channel. In the seco nd and third steps a general trend verified be fore is again observed: energy values are ba lanced by the respective densities of the cen ters involved. Long range hydrogen atom trans port is no longer observed as more stable com plex hydrogen aggregates appear.

The results obtained with the U₂ center \underline{mi} gration and aggregation into U_{2x} and H₂O centers added by the step transformation of U_{2x} into H_2O^- and, finally, the destruction of the $H_2O^$ center itself indicate that the highest"thermal"

energy barrier that the hydrogen atom in its ground state has to overcome to form an $H_20^$ defect must confine this atom in one of the 24 available interstitital sites next nearest neigh bours to the OH defect. The hydrogen atom thus trapped in one of these sites "feels" the polar ization effect due to the OH⁻ and forms the U_{2x} center but only with one transition of higher energy than the unperturbed U_2 center and with out splitting this transition as it happens with strongly perturbed U_2 center⁹. This polarization effect weakly affects the U_2 center resulting in the U_{2X} transition and confines the hydrogen atom into a more thermally stable place. We are presently extending this study to other alkali halide lattices and working with ESR techniques to confirm the present proposal of the U_{2x} center configurational model.

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